



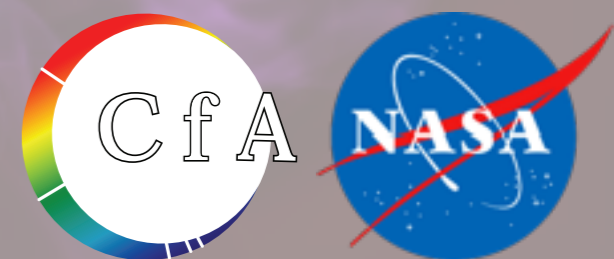
HOW OUTFLOWS AND RADIATIVE FEEDBACK LIMIT ACCRETION ONTO MASSIVE STARS

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Harvard-Smithsonian Center for Astrophysics

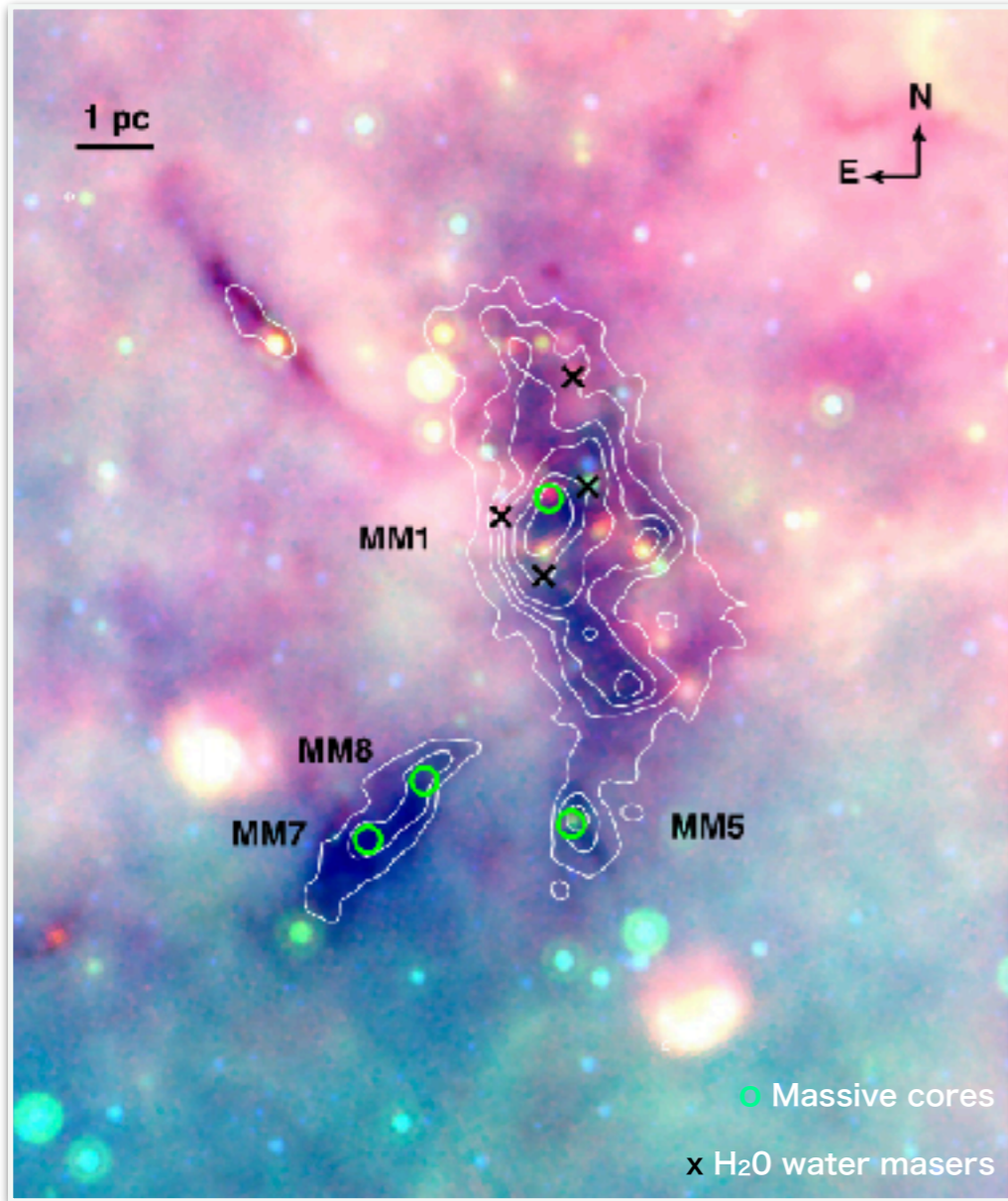
In collaboration with: Alyssa Goodman (CfA), Richard Klein (UCB/LLNL), Mark Krumholz (ANU), Aaron Lee (UT Austin), Chris McKee (UCB), Jeff Oishi (Bates College)



HARVARD-SMITHSONIAN
CENTER FOR ASTROPHYSICS

Massive star formation is [likely] a scaled up version of low-mass star formation

Infrared dark cloud (IRDC) G28.53



Lu+2015

- IRDCs can fragment into dense, massive clumps which then fragment into massive pre-stellar cores.
- Massive pre-stellar cores are supported by turbulent pressure
$$P_{\text{Turb}} \gg P_{\text{Th}}$$
- Observations suggest massive cores have $\alpha_{\text{vir}} \lesssim 1$

$$\alpha_{\text{vir}} = \frac{2E_{\text{KE}}}{E_{\text{G}}} = \frac{5\sigma^2 R_c}{GM_c}$$

Isotropic accretion leads to the radiation pressure barrier problem in massive star formation

Formation of massive stars is a competition between gravity and (direct+indirect) radiation pressure

Gravitational Force:

$$f_{\text{grav}}(r) = \frac{GM_{\star}\Sigma}{r^2}$$

$$\Sigma(r) = \int_0^r \rho(r') dr'$$

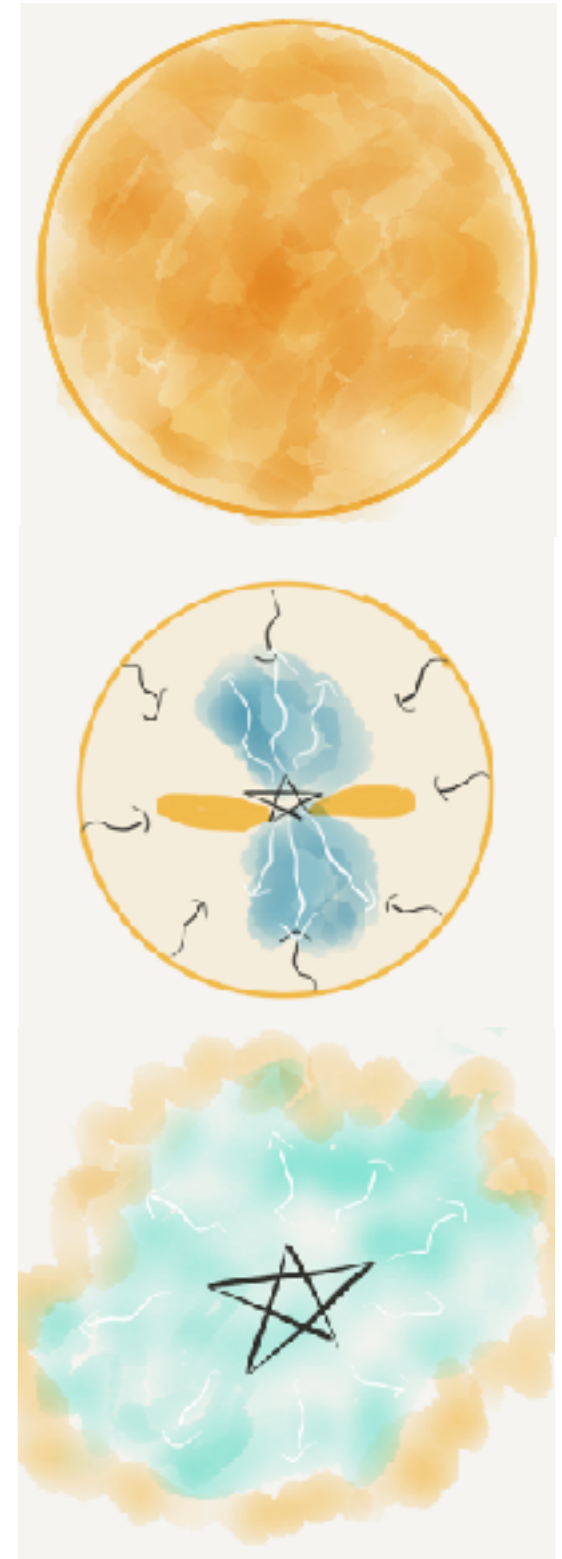
Radiative Force:

$$f_{\text{rad}} = \frac{L_{\star}}{4\pi r^2 c} (1 + f_{\text{trap}})$$

$$L_{\star} \propto M_{\star}^3$$

$$f_{\text{edd}} = 7.7 \times 10^{-5} (1 + f_{\text{trap}}) \left(\frac{L_{\star}}{M_{\star}} \right)_{\odot} \left(\frac{\Sigma}{1 \text{ g cm}^{-2}} \right)^{-1}$$

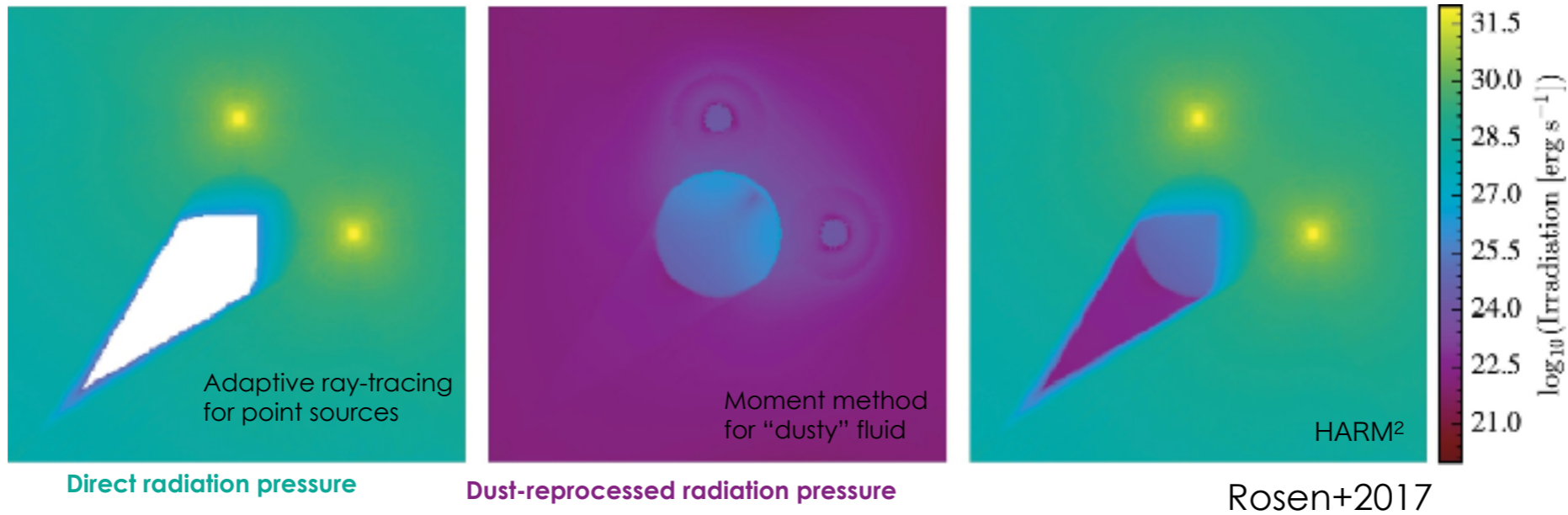
Radiation halts isotropic accretion when $f_{\text{edd}} \gtrsim 1$
for $M_{\star} \gtrsim 20 M_{\odot}$



Modeling massive star formation
requires **multi**-dimensional
radiation-hydrodynamic
simulations

Modeling radiation pressure in (massive) star formation simulations

Hybrid Adaptive Ray-Moment Method (**HARM²**):



Absorption of (multi-frequency) stellar radiation field:

Radiative Transfer
Equation along ray:

$$\frac{\partial L_{\text{ray},j}}{\partial r} = -\kappa_j \rho L_{\text{ray},j},$$

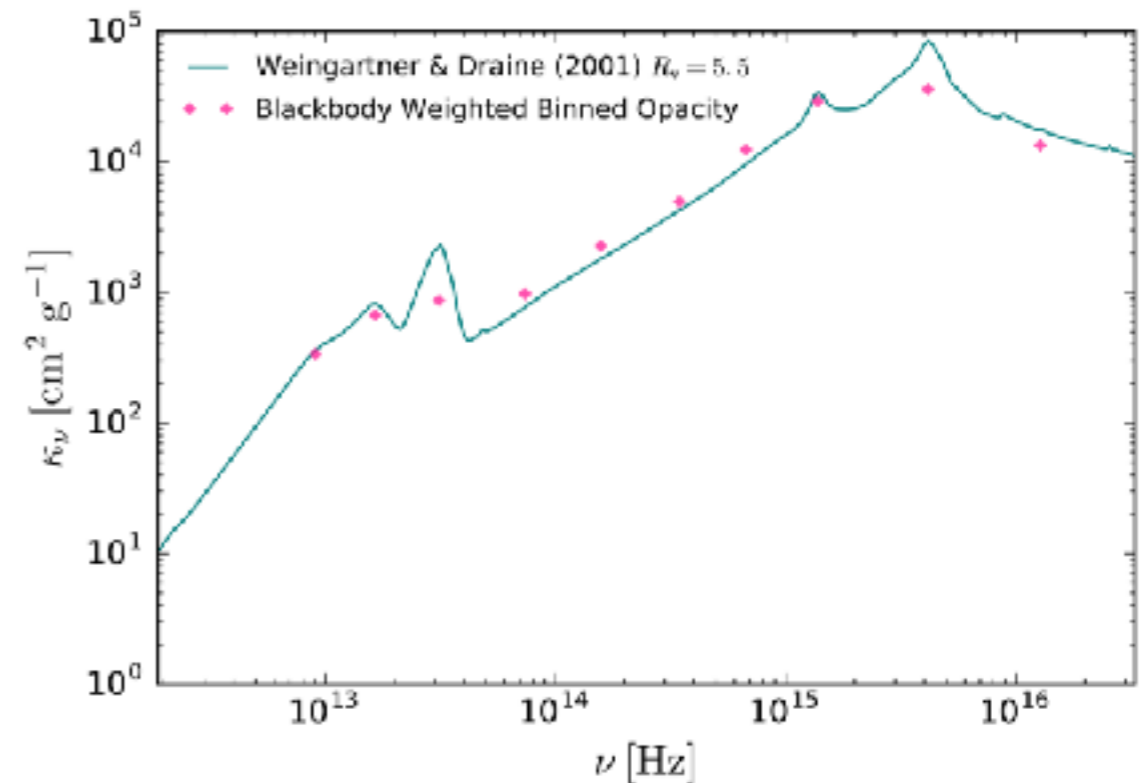
Luminosity absorbed
($\tau_j = \rho_d \kappa_j dl$):

$$dL_{\text{ray},j} = L_{\text{ray},j} (1 - e^{-\tau_j})$$

Energy and momentum deposition:

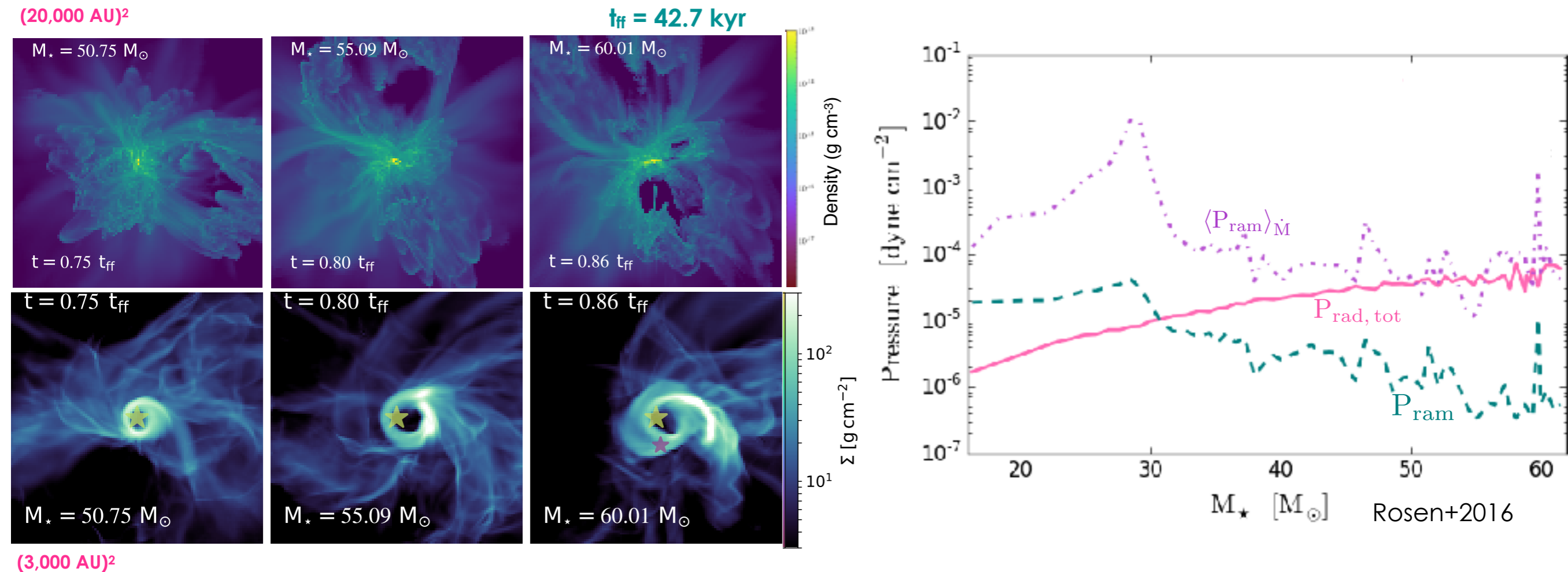
$$\dot{\epsilon}_{\text{rad,ray}} = \sum_{j=1}^{N_\nu} dL_{\text{ray},j}$$

$$\dot{\mathbf{p}}_{\text{rad,ray}} = \sum_{j=1}^{N_\nu} \frac{dL_{\text{ray},j}}{c} \mathbf{n}.$$



Rosen+2016

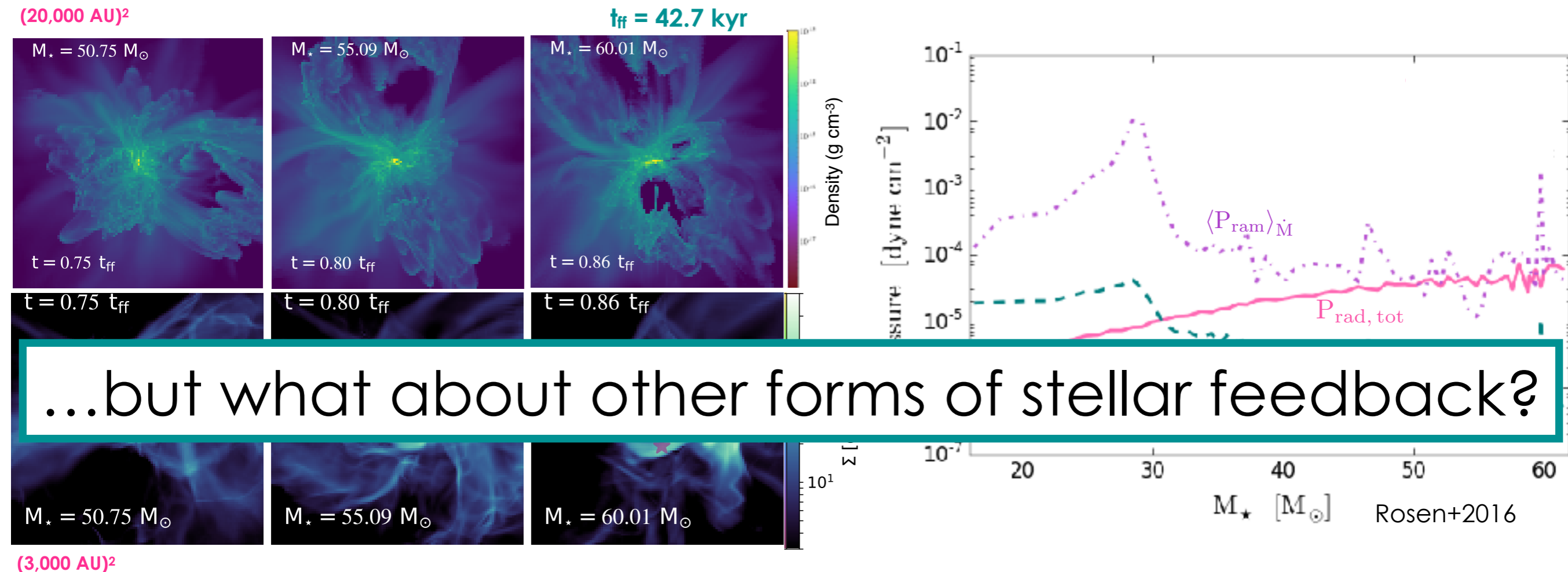
Overcoming the radiation pressure barrier



Mass delivered to star via infalling dense filaments, radiative Rayleigh Taylor (RT) instabilities, and disk accretion.

High accretion rates and infalling filaments provide sufficient ram pressure to overcome radiation pressure.

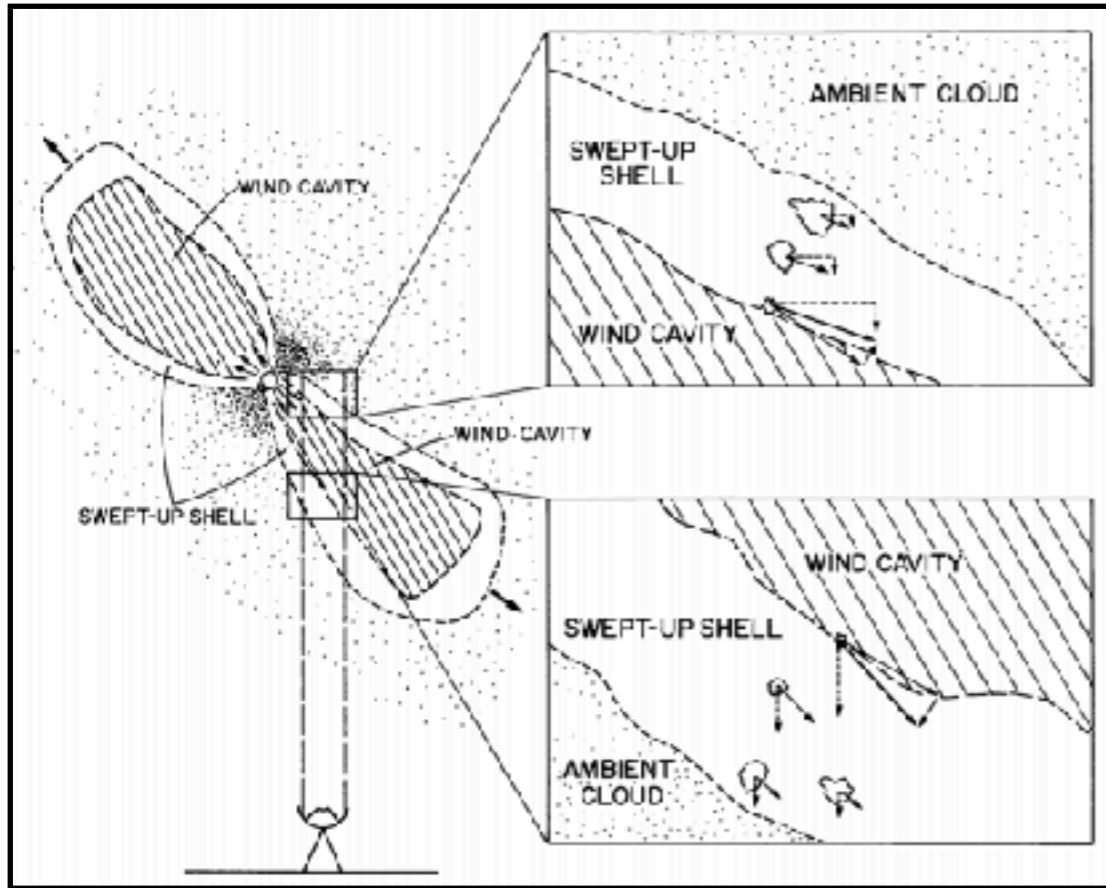
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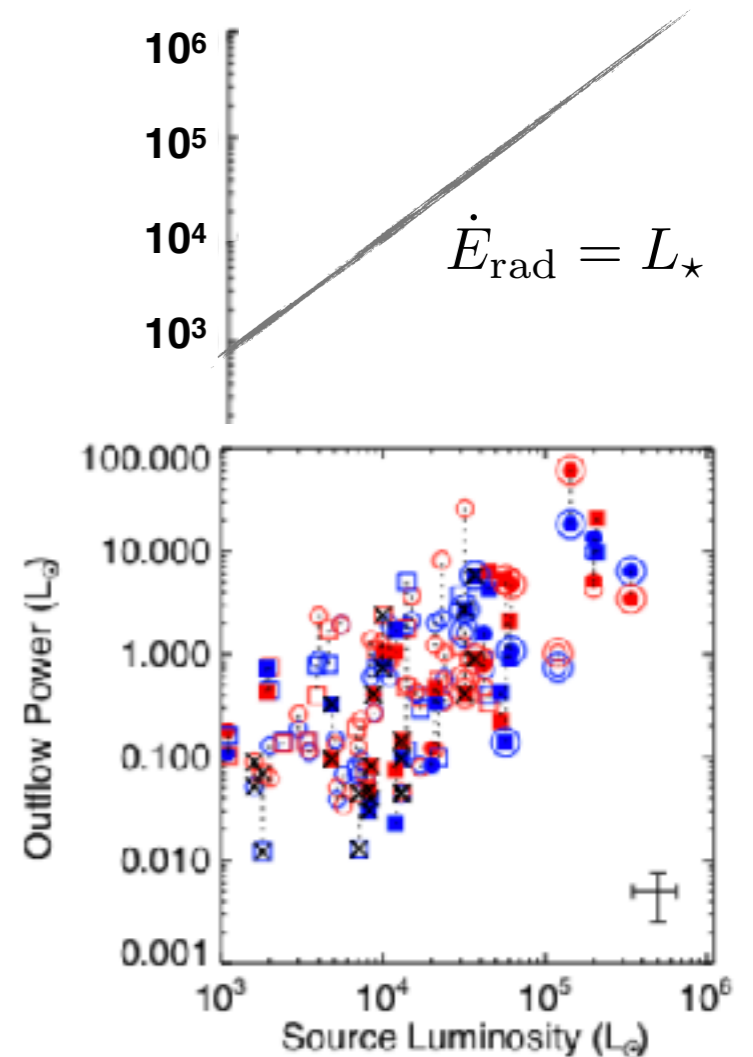
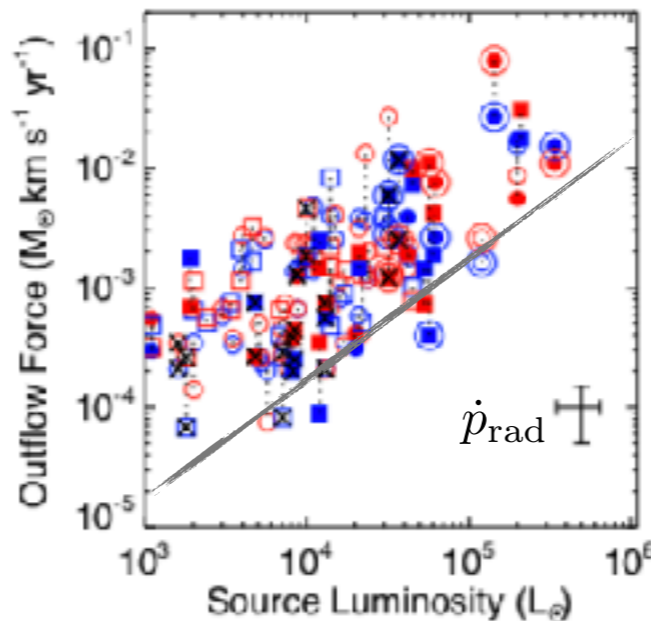
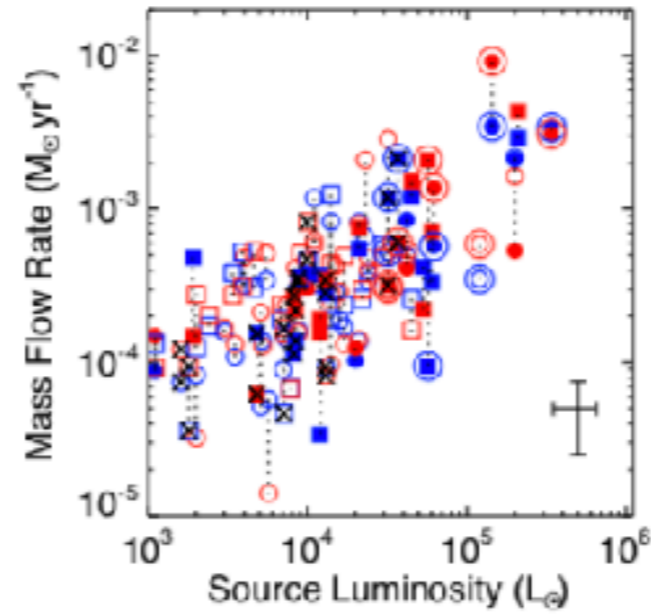
High accretion rates and infalling filaments provide sufficient ram pressure to overcome radiation pressure.

Collimated bipolar outflows are ubiquitous in (low-mass and) high-mass star formation



Lada 1985, ARAA

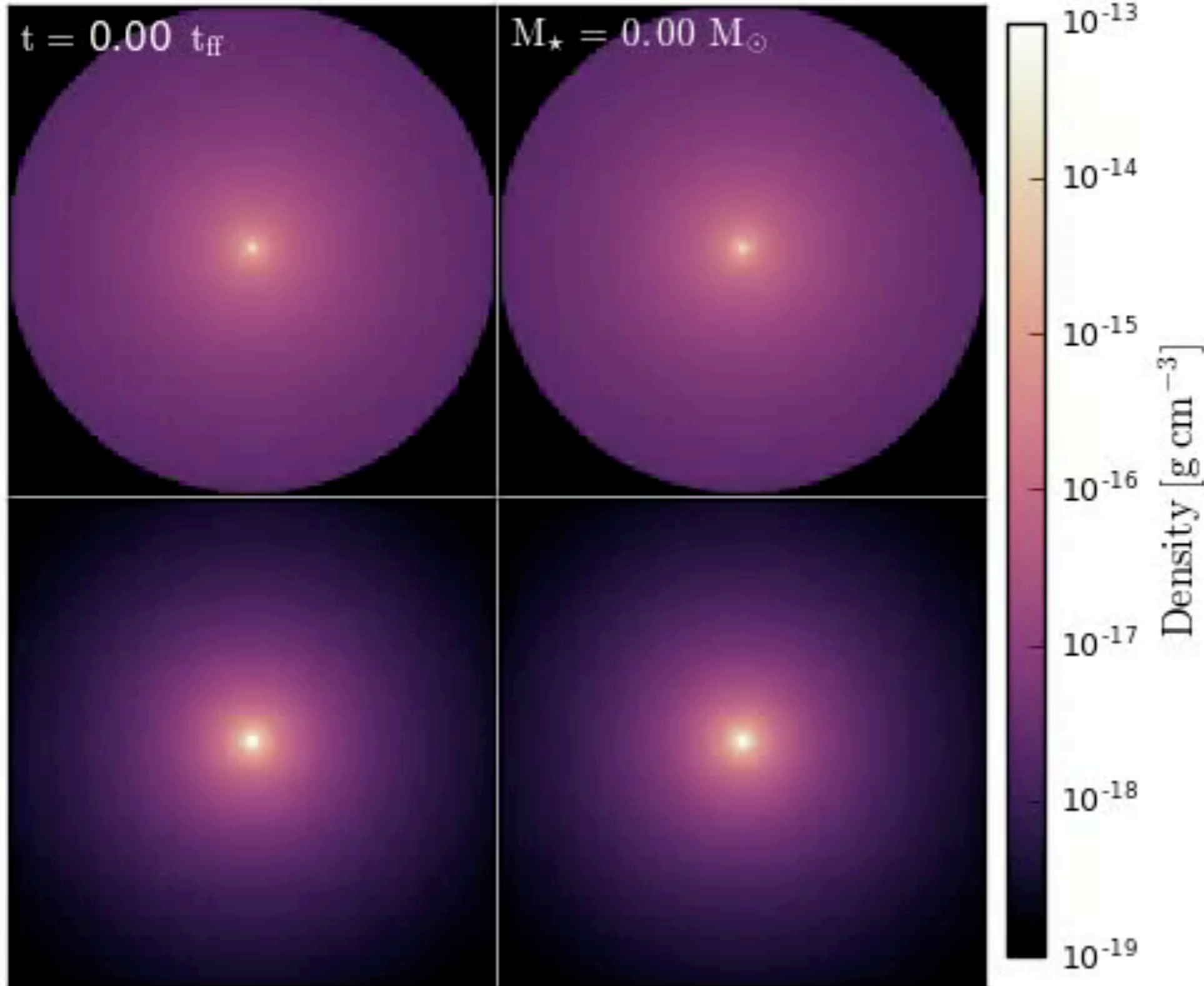
$$v_{\text{jet}} \sim \sqrt{\frac{GM_{\star}}{R_{\text{star}}}} \sim 100 \text{ km s}^{-1}$$



$$\dot{p}_{\text{rad}} = \frac{L_{\star}}{c} = 2 \times 10^{-8} \left(\frac{L_{\text{star}}}{L_{\odot}} \right) M_{\odot} \text{ km s}^{-1} \text{ yr}^{-1} \quad \text{Maud+2015}$$

Powerful jets from accreting stars can drive wide angle molecular outflows from star-forming cores and eject core material

Massive star formation with radiative and outflow feedback



Initial Conditions:

$$M_{\text{core}} = 150 M_{\odot}$$

$$R_{\text{core}} = 0.1 \text{ pc}$$

$$\rho(r) \propto r^{-3/2}$$

$$\sigma_{1D} = 1.2 \text{ km s}^{-1}$$

$$\alpha_{\text{vir}} \sim 1$$

$$\Delta x_{\text{min}} = 20 \text{ AU}$$

$$t_{\text{ff}} = 42,710 \text{ yrs}$$

$$p_{\text{OF}} = \dot{M}_{\text{OF}} v_{\text{OF}}$$

$$\dot{M}_{\text{OF}} = 0.21 \times \dot{M}_{\text{acc}}$$

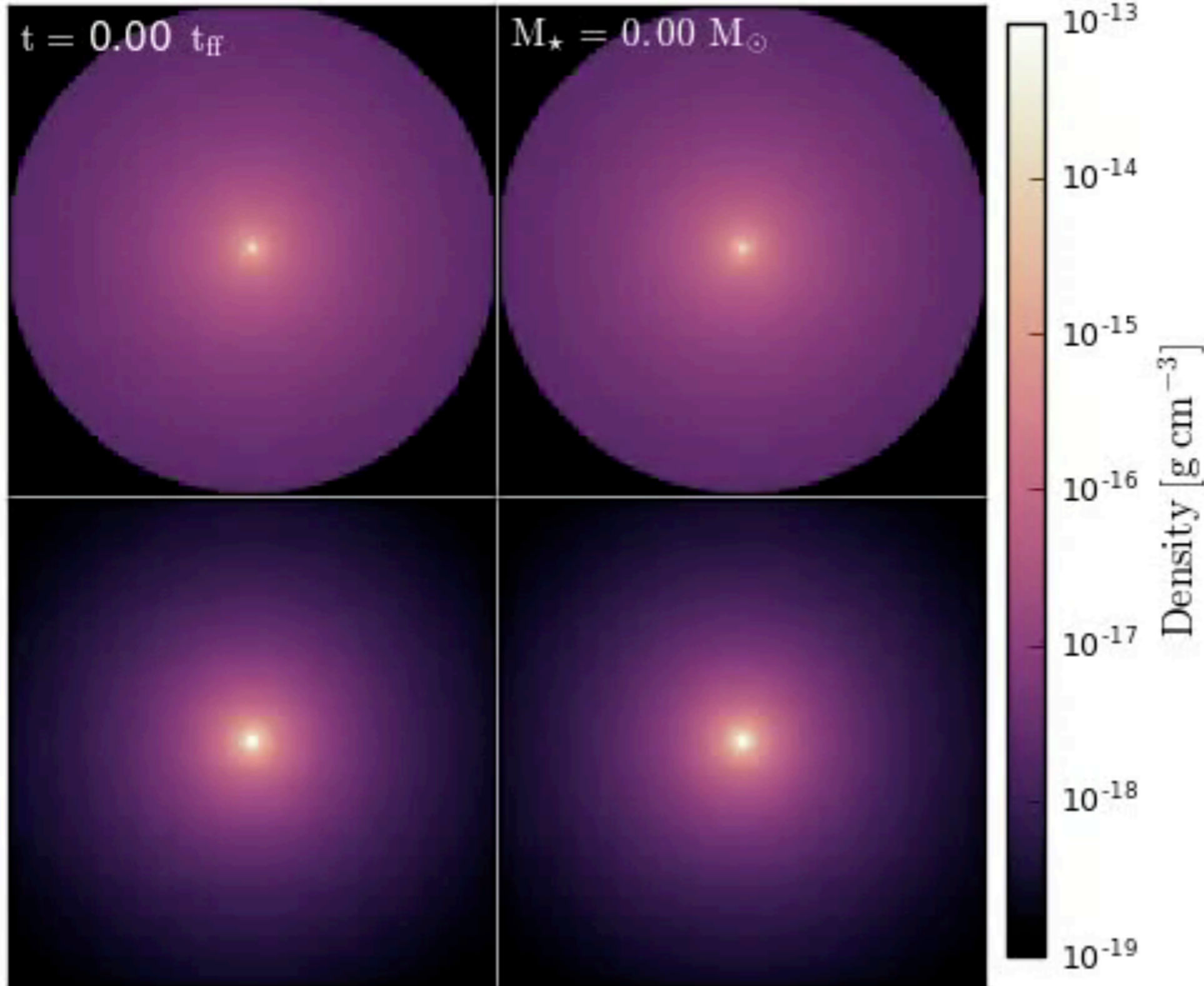
$$v_{\text{OF}} = 0.3 \times v_{\text{esc}}$$

Rosen+(in prep)

Top panel: (40,000 AU x 40,000 AU)

Bottom panel: (8,000 AU x 8,000 AU)

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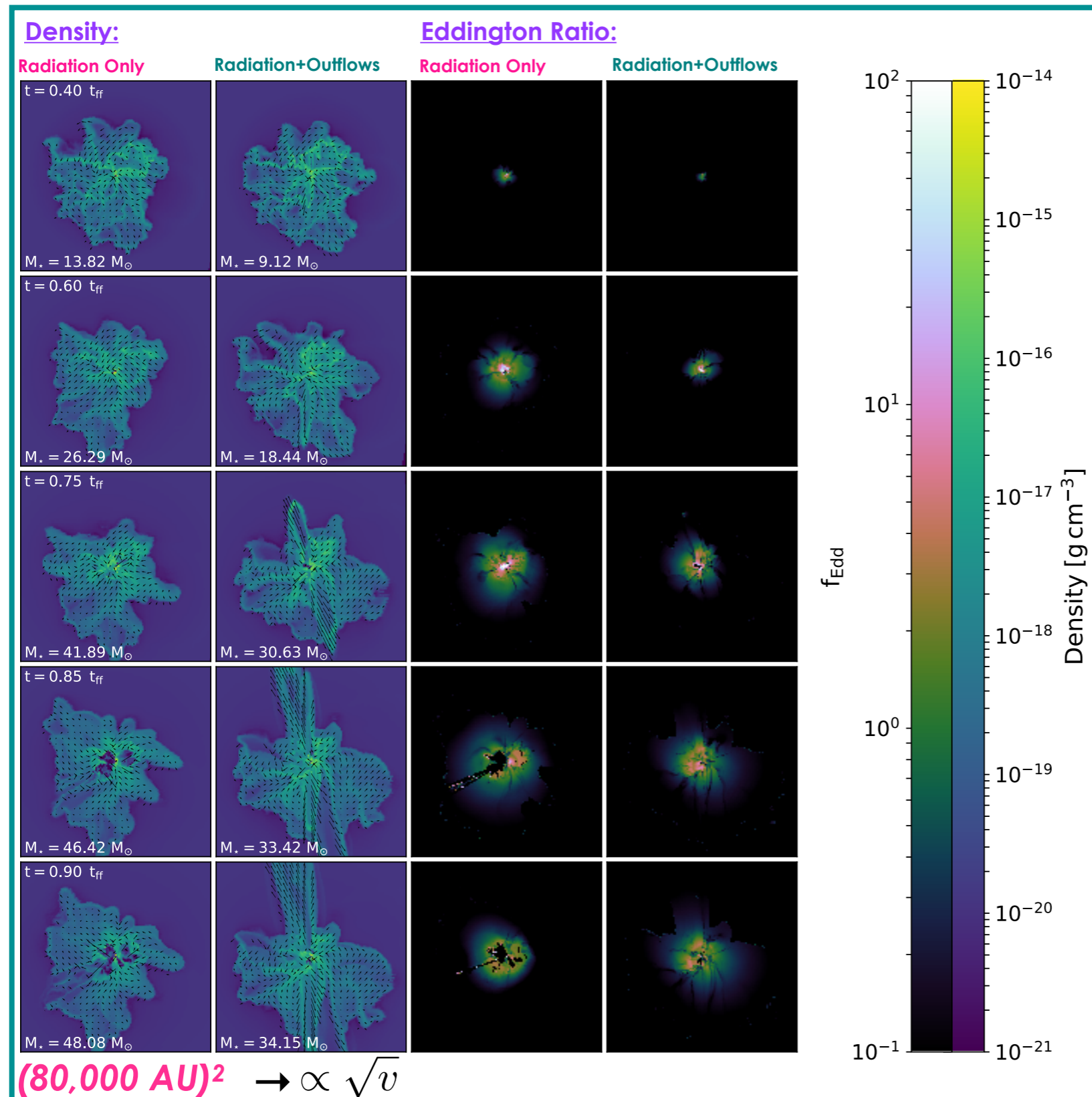
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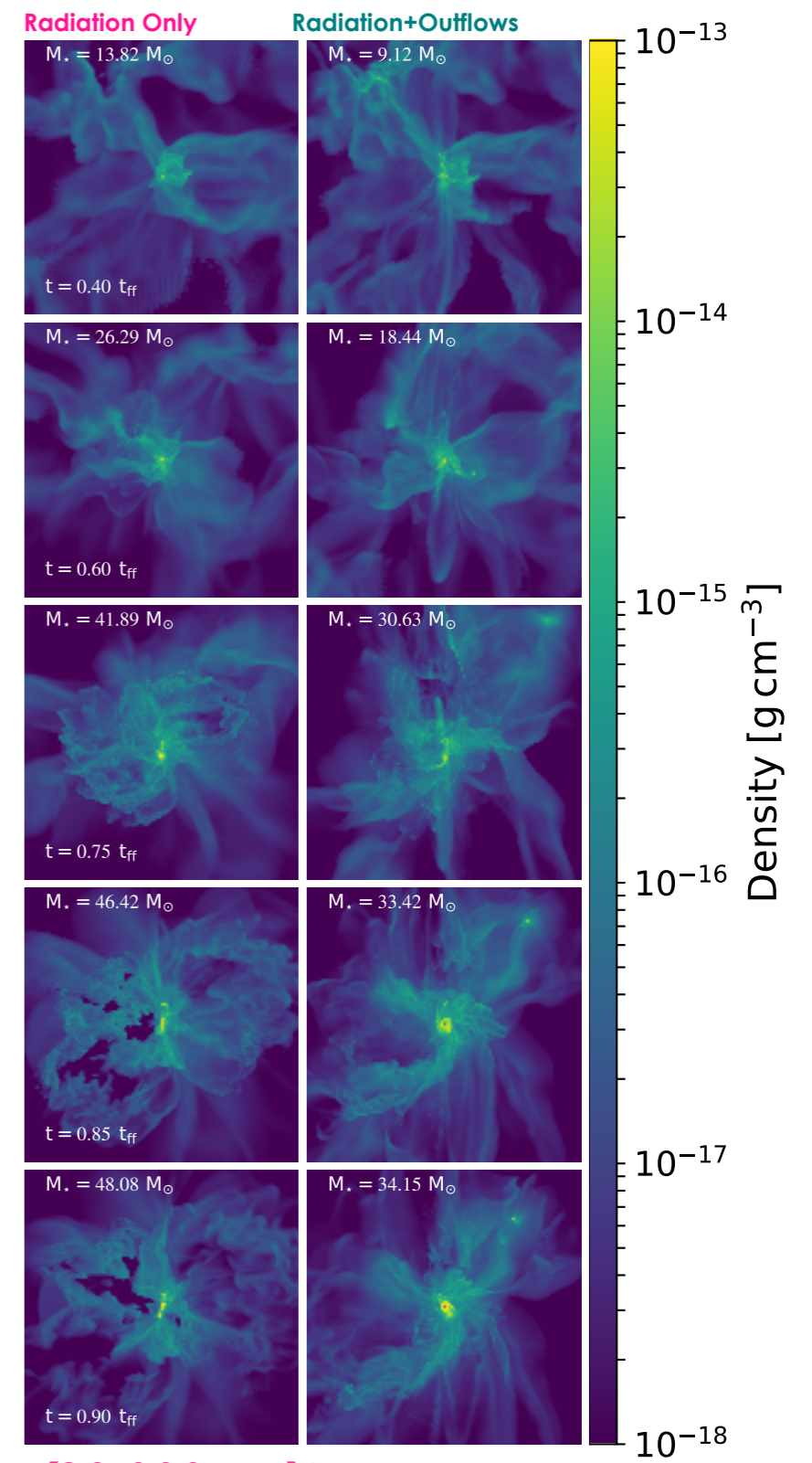
Top panel: (40,000 AU x 40,000 AU)

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Outflows punch holes in ISM along the star's polar directions **allowing radiation to escape**, thereby reducing the development of RT instabilities.

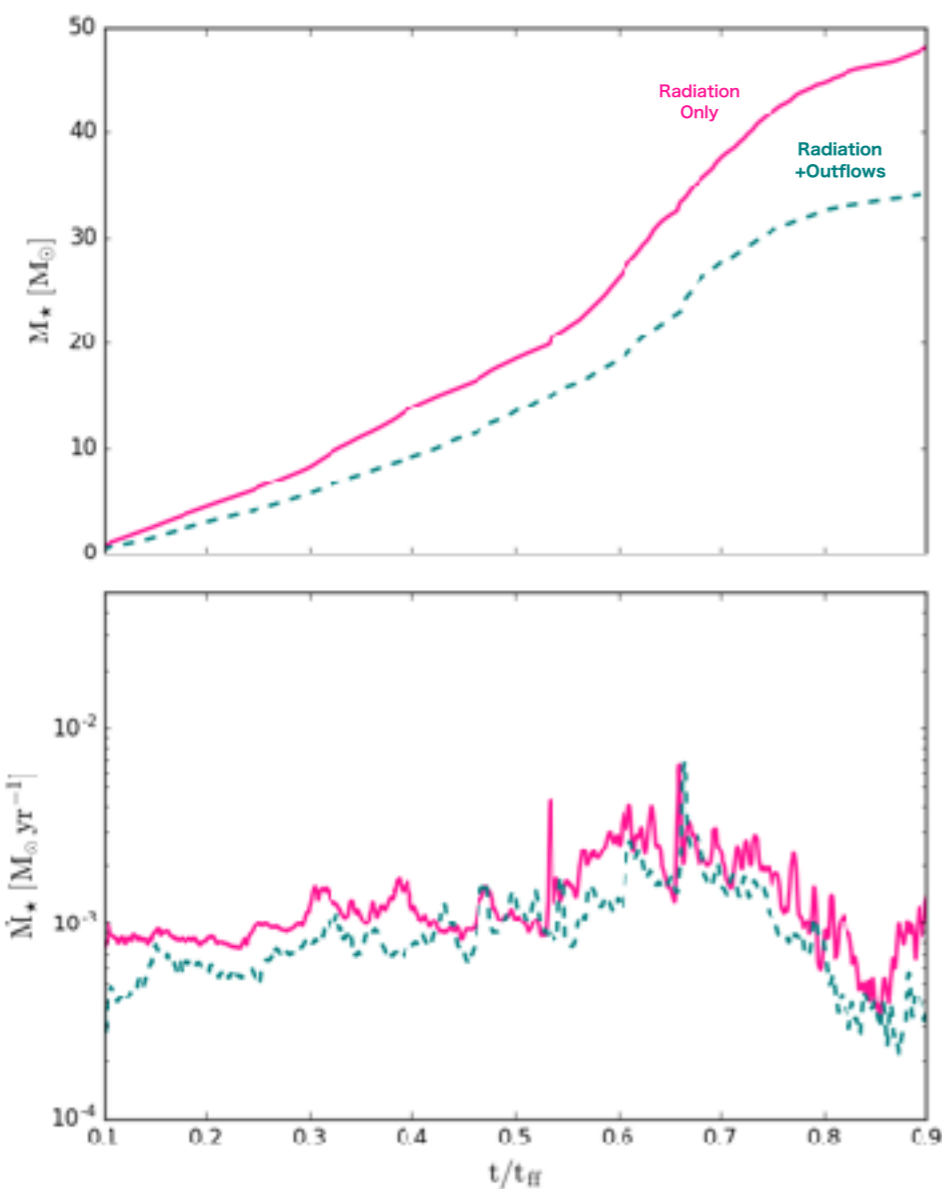


Thin Density Projections:



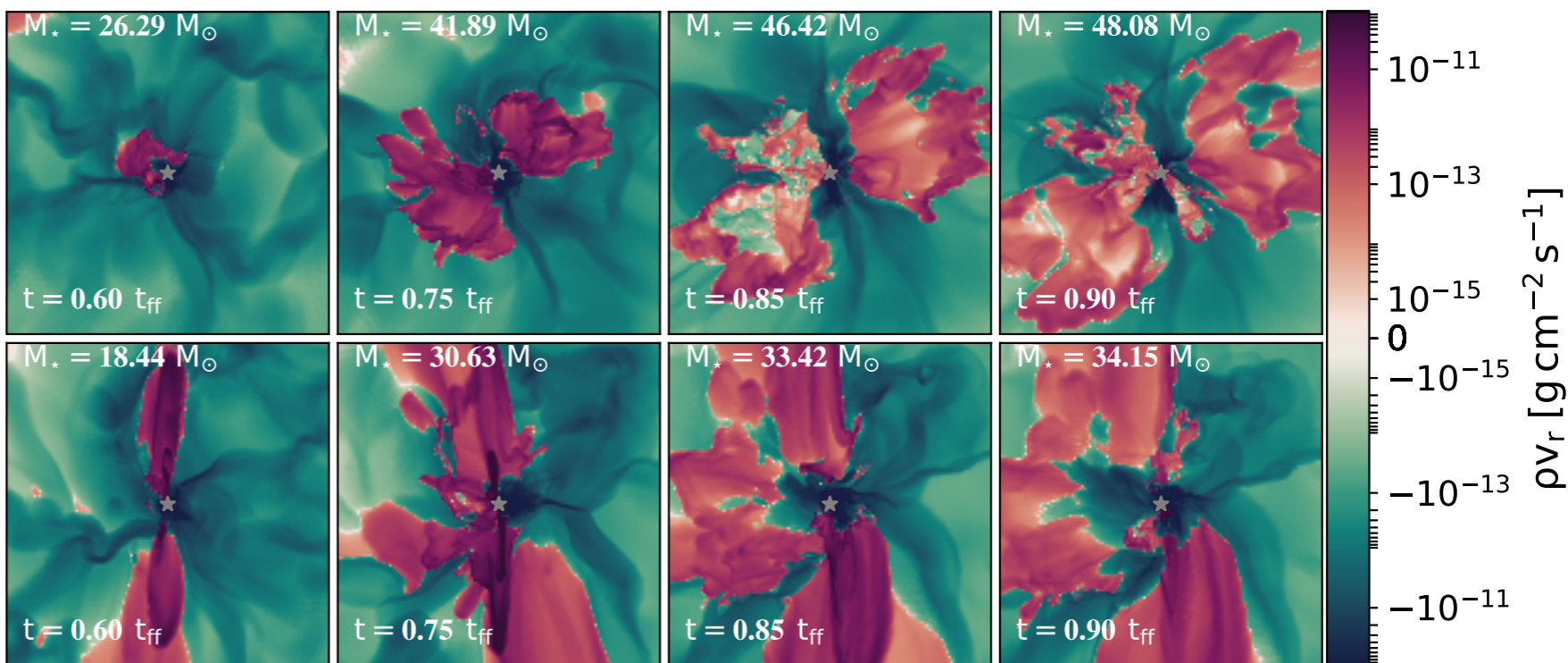
Rosen+(in prep)

Outflows+radiation pressure **efficient at ejecting material** away from the star than radiation pressure alone.



Rosen+(in prep)

Radiation Only



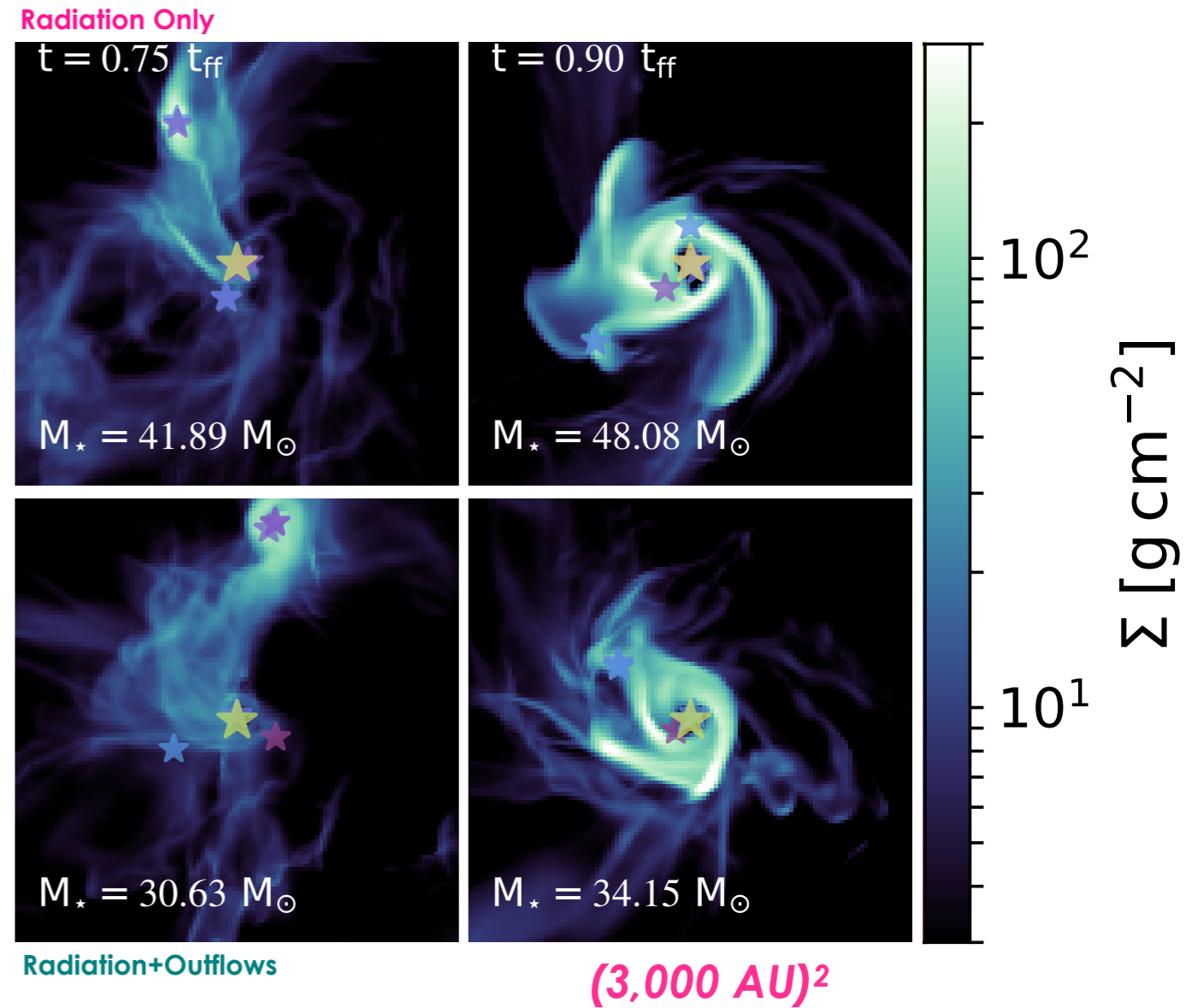
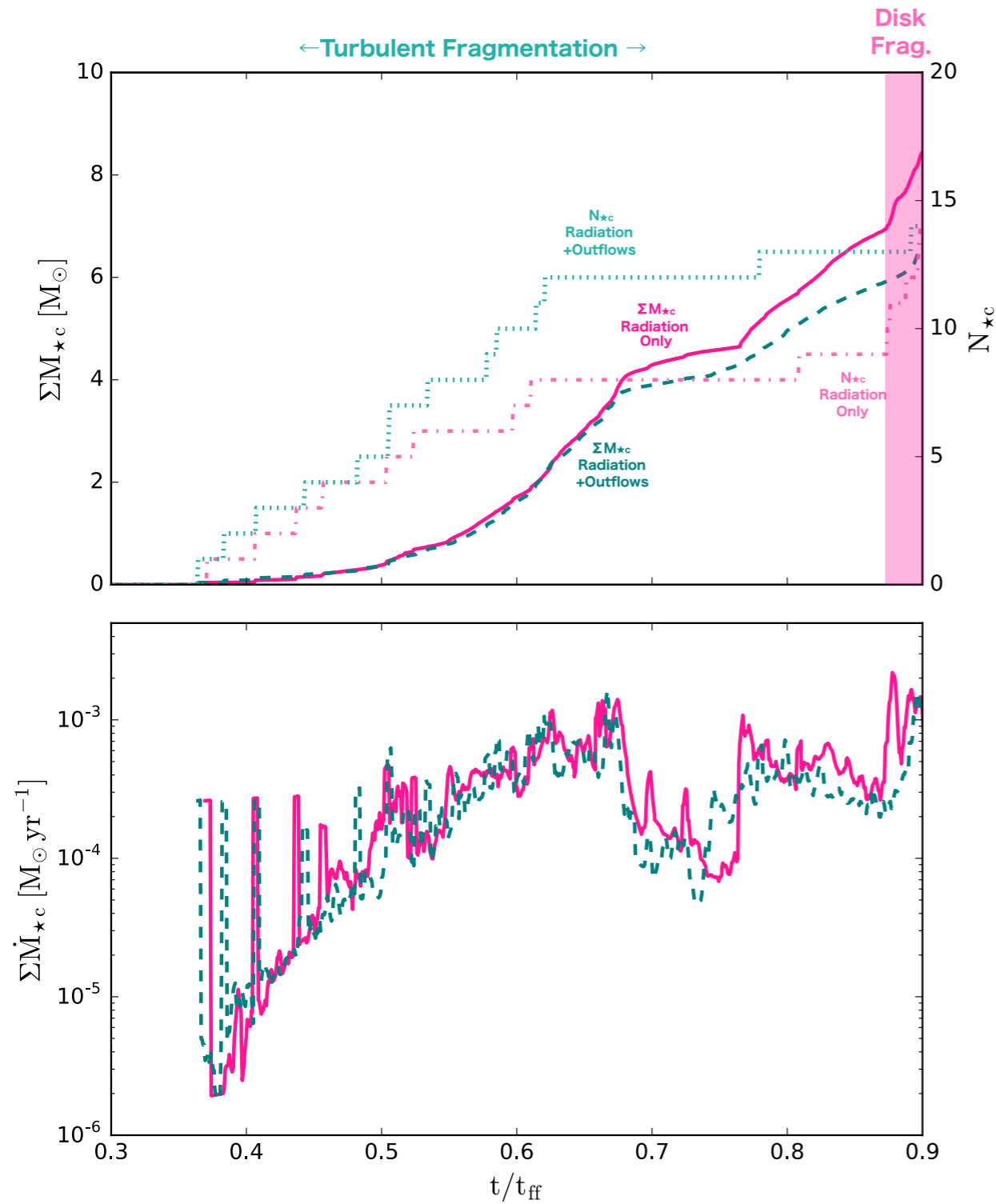
Radiation+Outflows

(20,000 AU)²

21% of accreted mass ejected by outflows

→ **only ~8% difference!**

Disks are crucial to massive star formation, especially at late times.

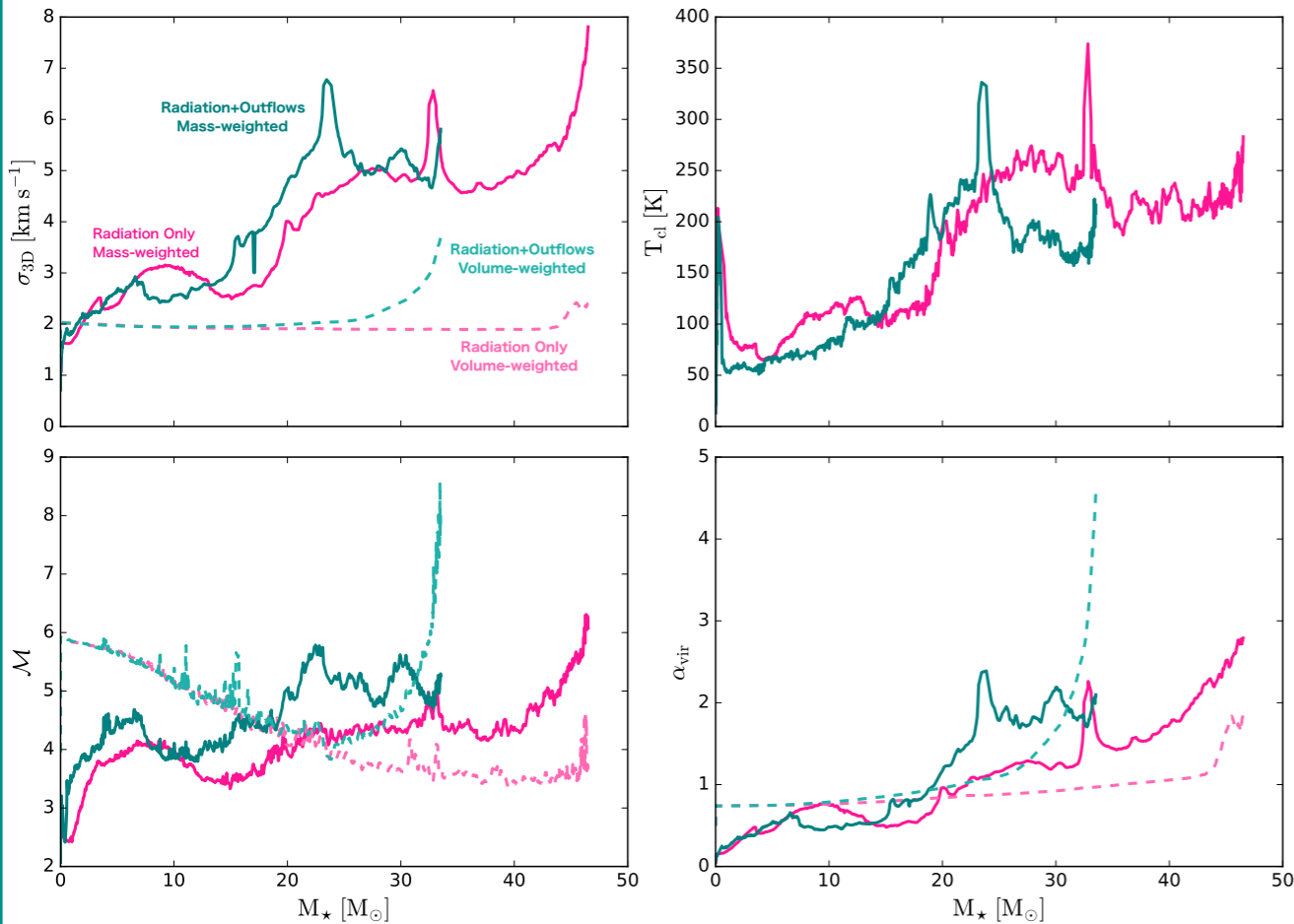


Rosen+(in prep)

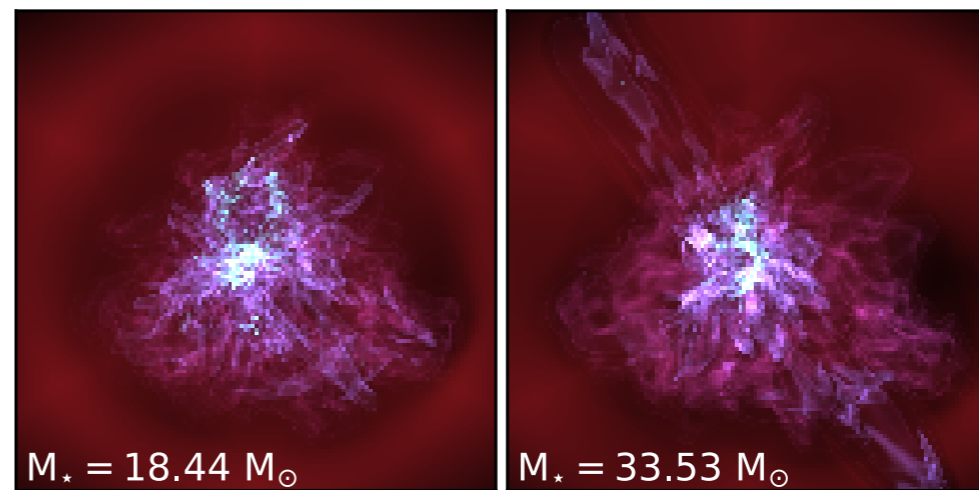
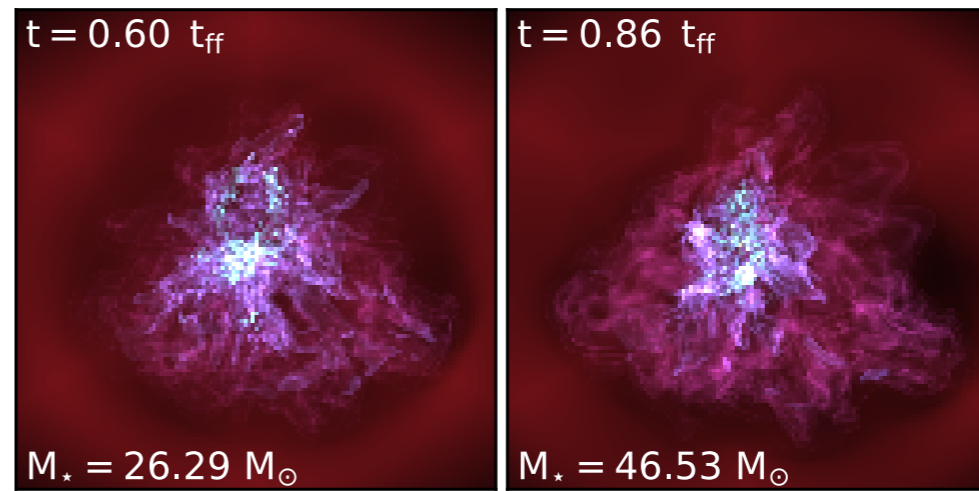
Companions formed via turbulent fragmentation at early times, disk fragmentation at late times

Outflows drive out entrained gas, eventually unbinding the core

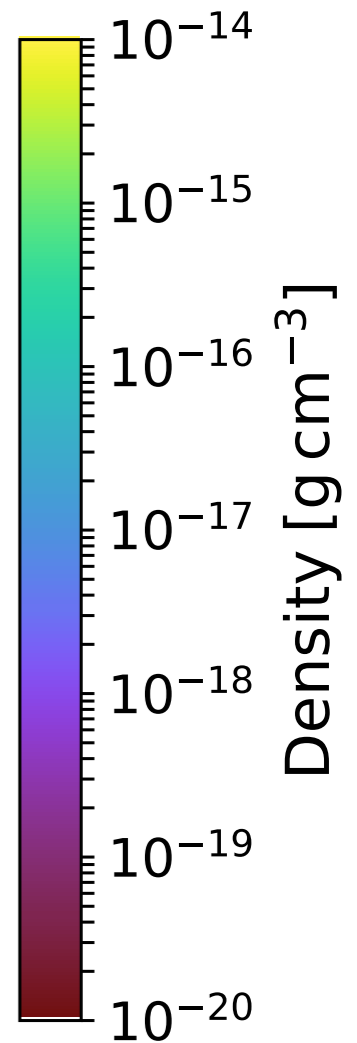
Core properties vs. M_\star



Radiation Only



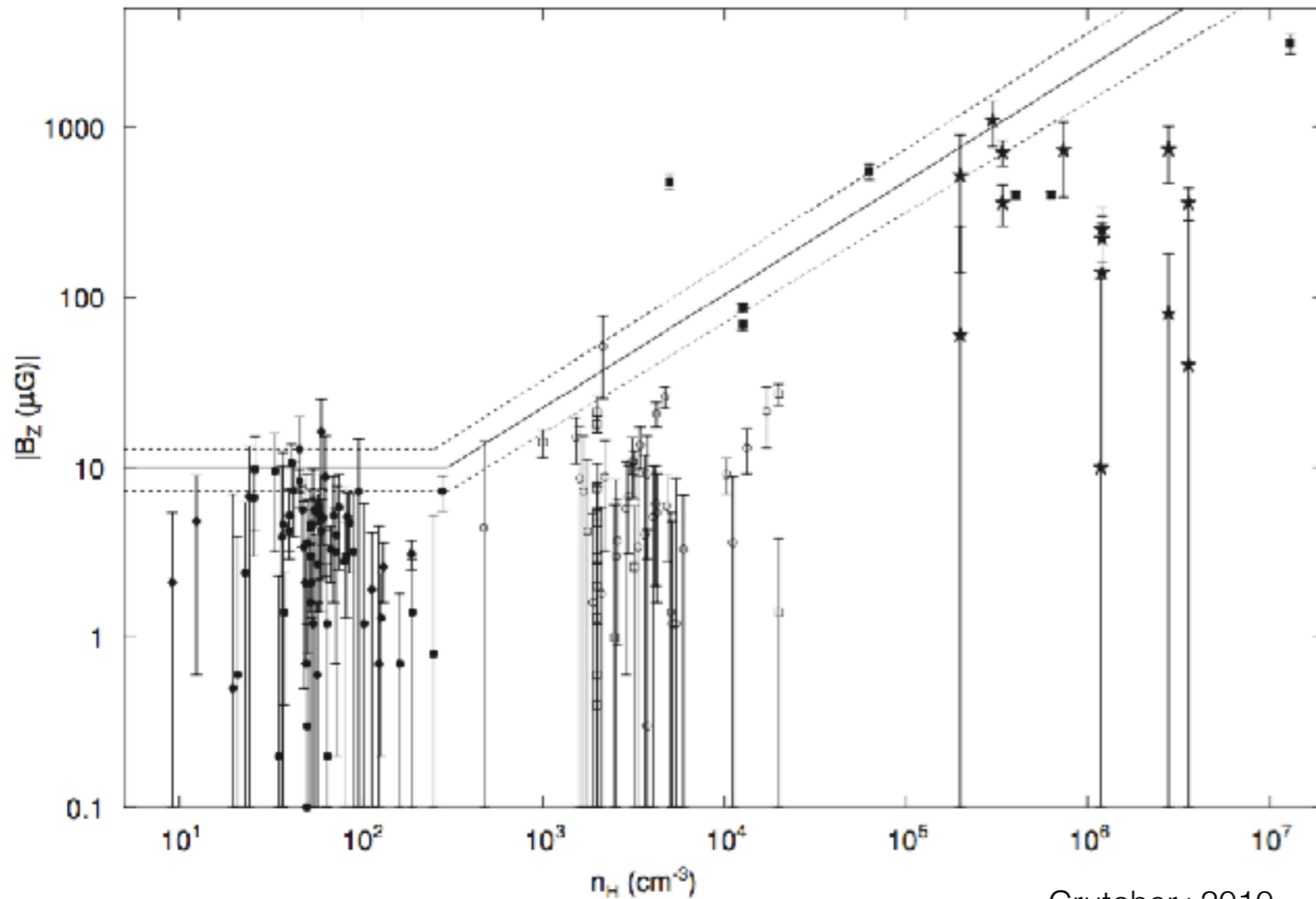
Radiation+Outflows



Rosen+(in prep)

Feedback from outflows allows radiation to escape, thereby reducing radiative heating.

...**BUT WAIT!** What about magnetic fields?



$$\mu_{\Phi} = \frac{M}{M_{\Phi}} \simeq \frac{2\pi\sqrt{GM}}{\pi B^2}$$

Supercritical

$$\mu_{\Phi} > 1$$

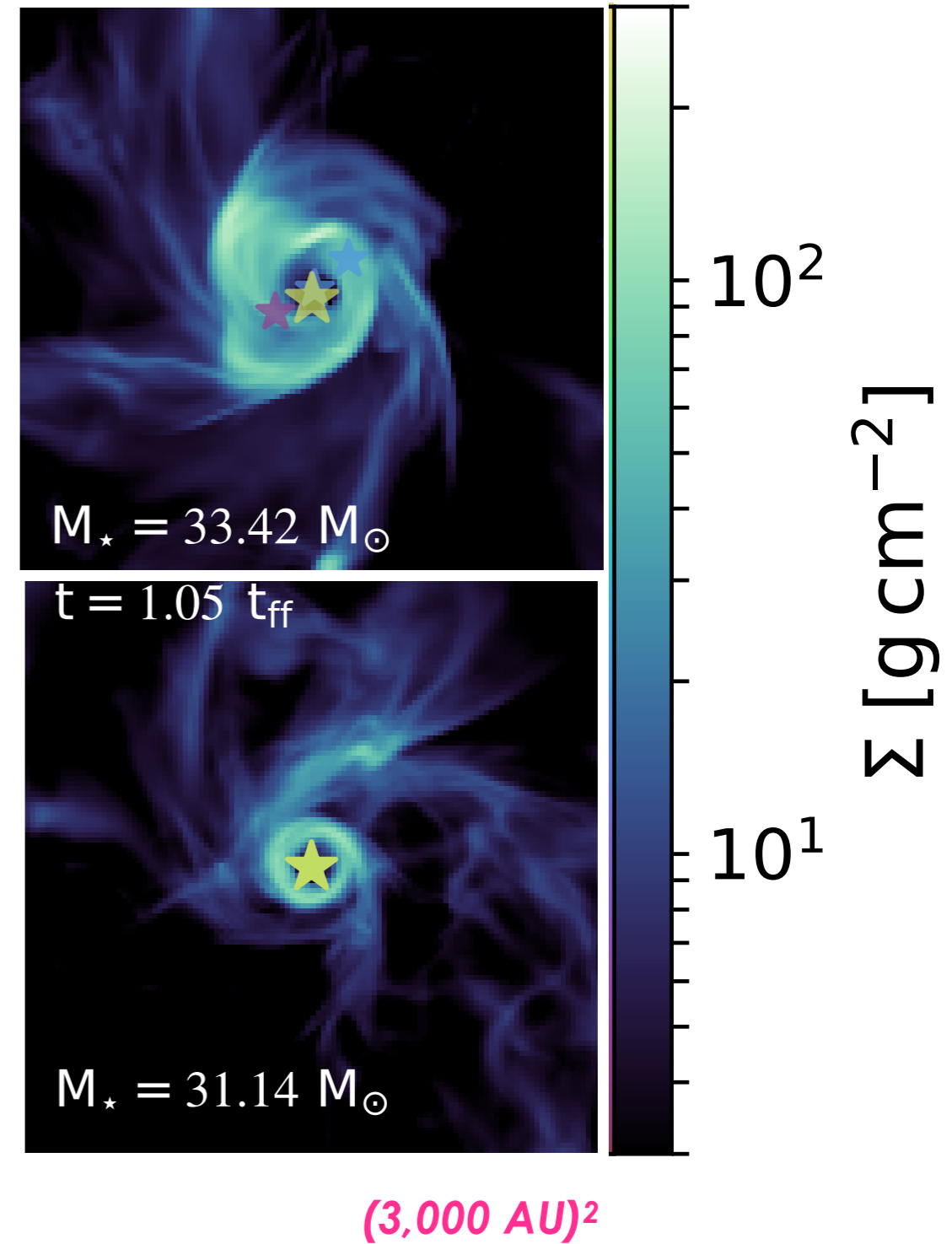
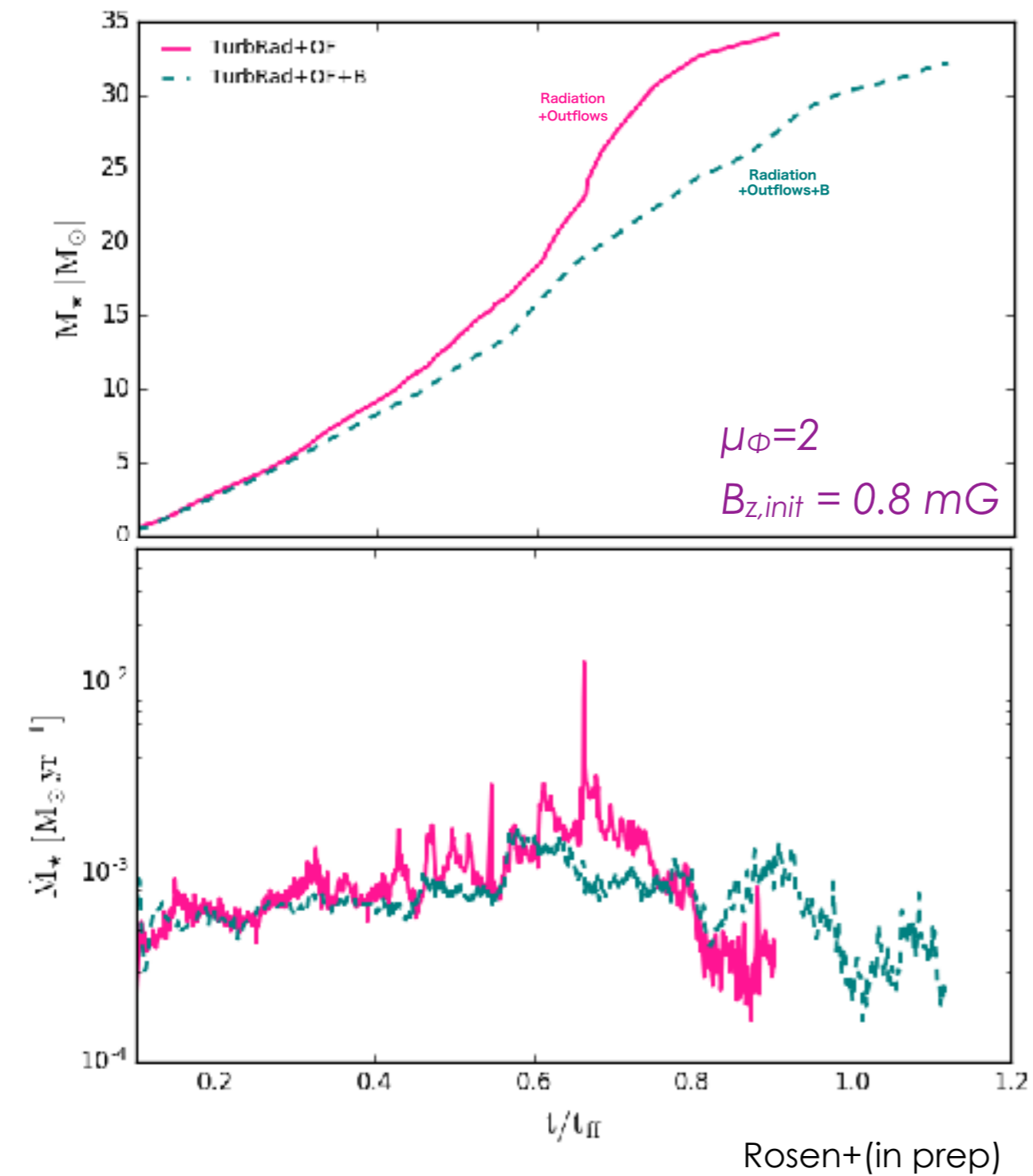
Subcritical

$$\mu_{\Phi} < 1$$

Observations suggest that dense molecular gas has $\mu_{\Phi} \sim 2$ (supercritical).

Magnetic pressure will **slow down collapse** and **reduce fragmentation**.

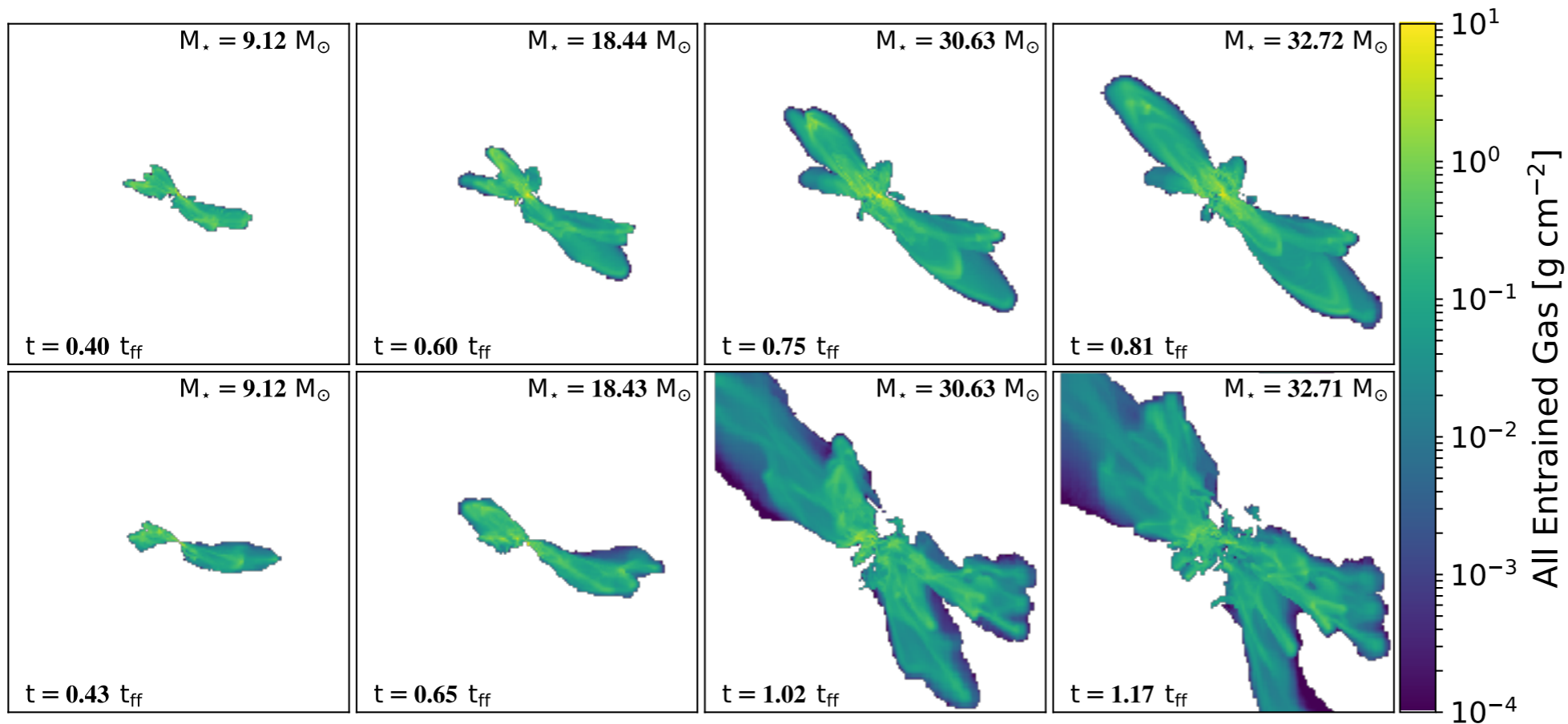
Magnetic braking **removes angular momentum** resulting in a smaller disk. Fragmentation is **highly suppressed**.



Inclusion of magnetic fields **reduces final stellar mass** by $\sim 20\%$ @ $t = 0.9 t_{ff}$

Entrained molecular outflows are collimated, but have **wider opening angles** when magnetic fields are included

Radiation+Outflows



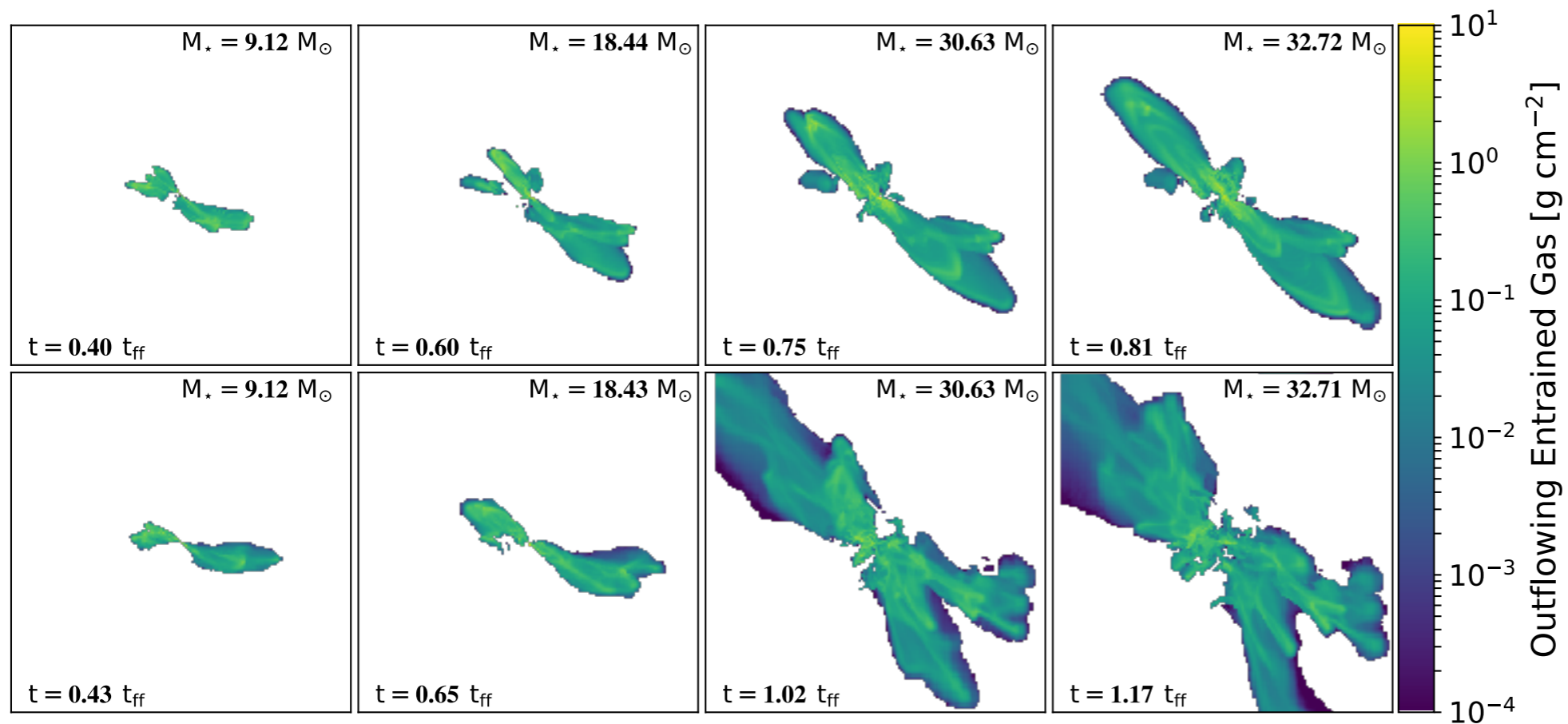
$$\rho_{\text{OF}} / \rho \gtrsim 5\%$$

Radiation+Outflows+B

Rosen+(in prep)

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Radiation+Outflows



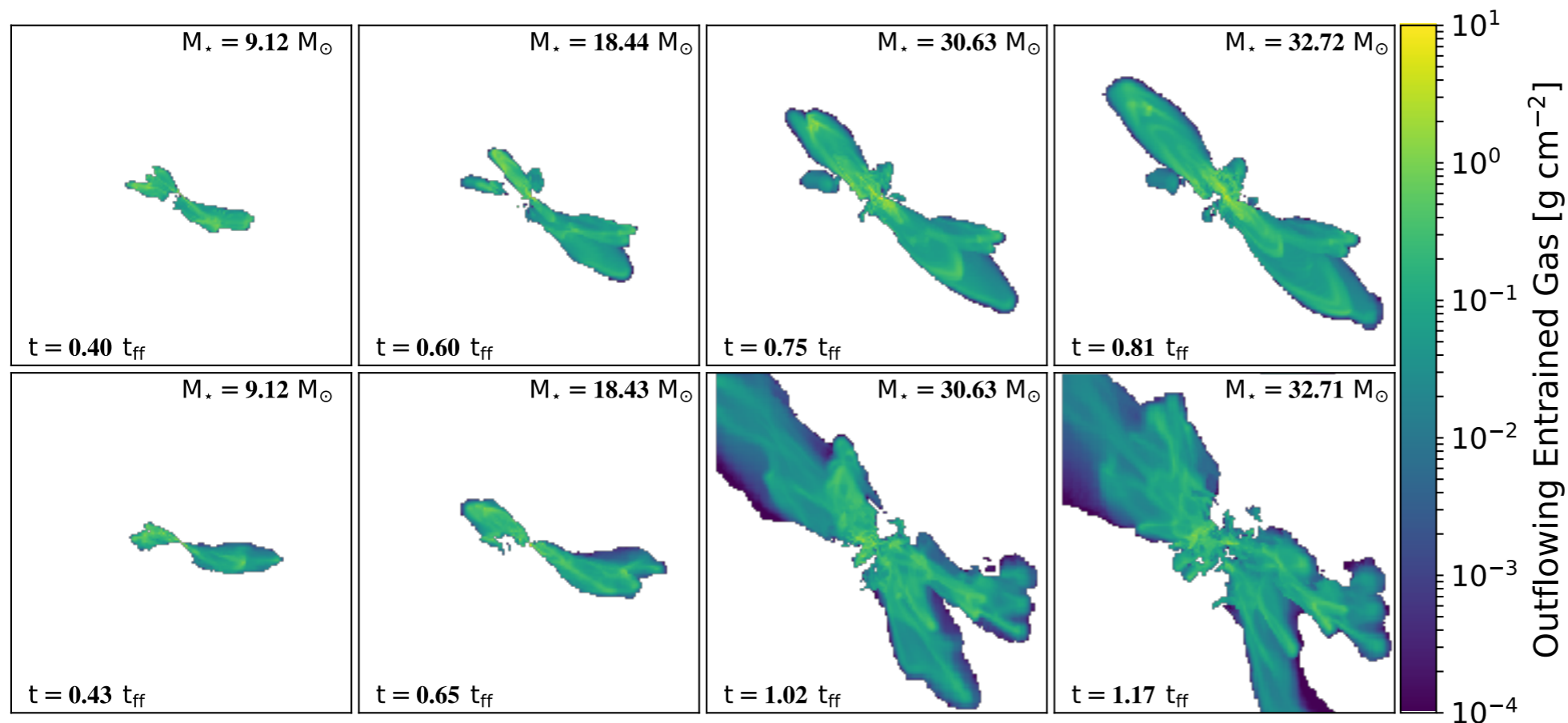
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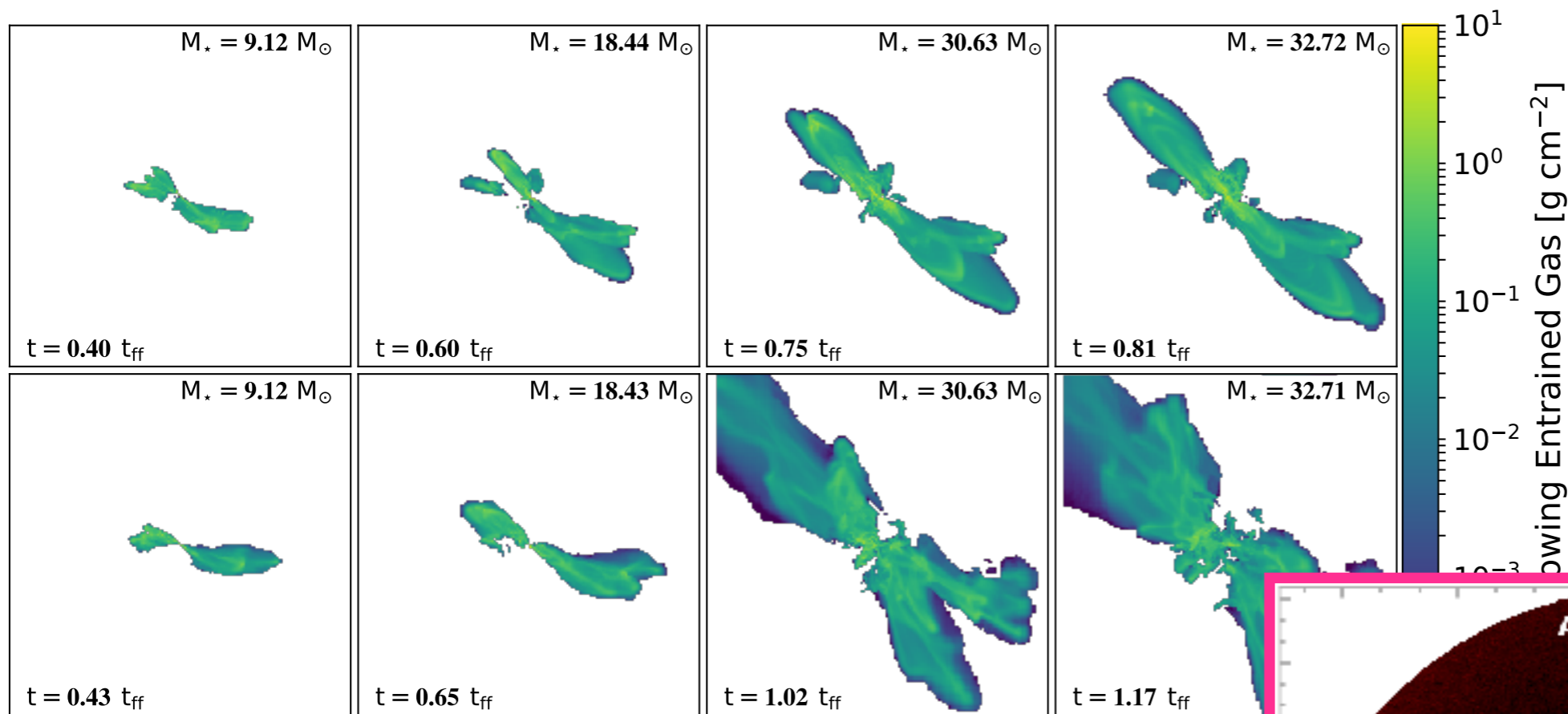
Radiation+Outflows+B

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...but how does this compare to observations?

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Radiation+Outflows



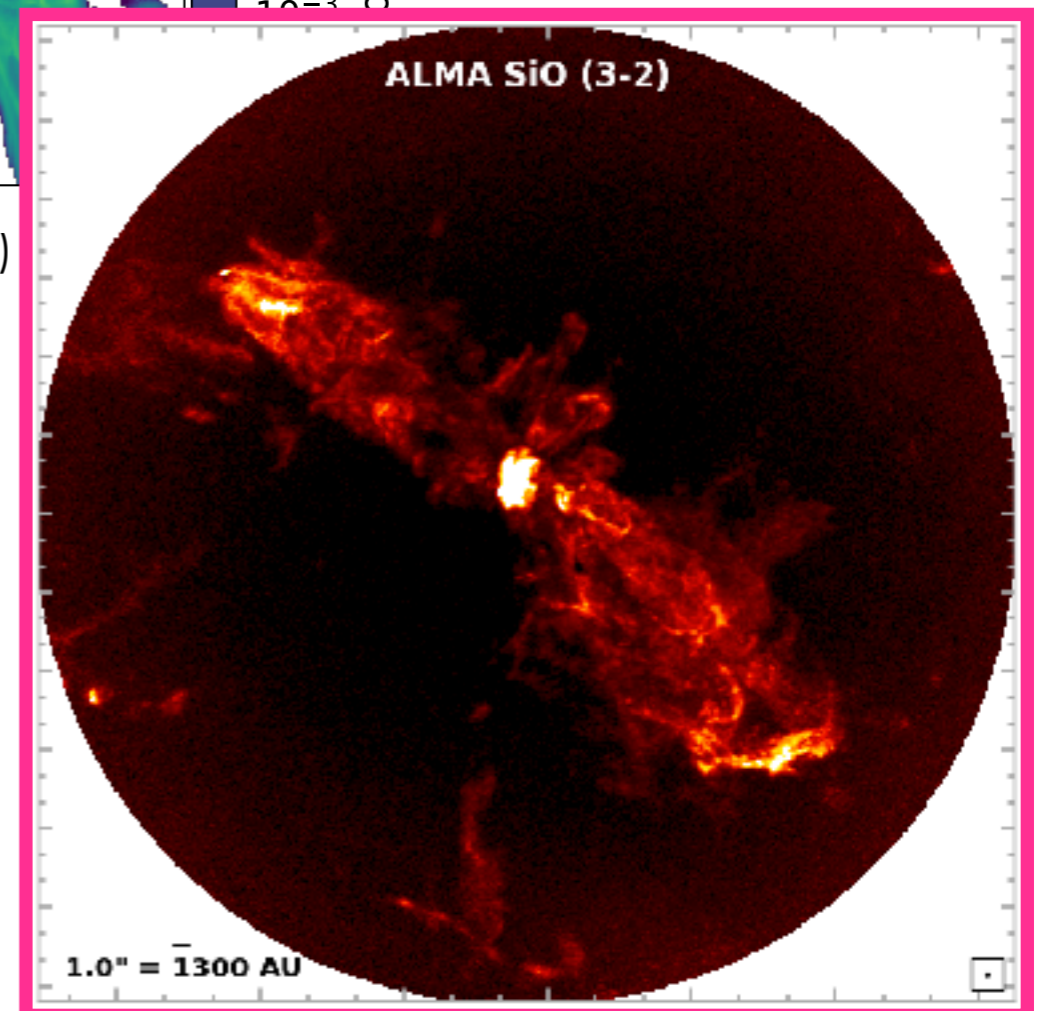
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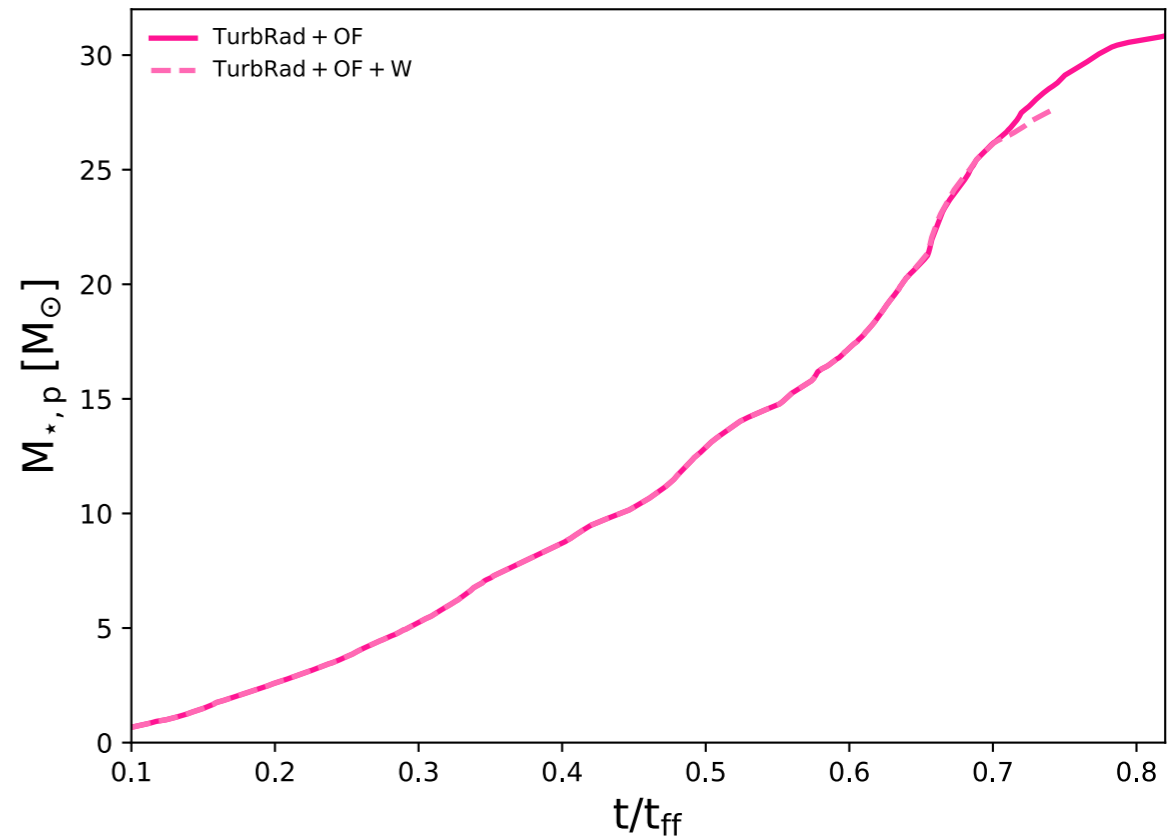
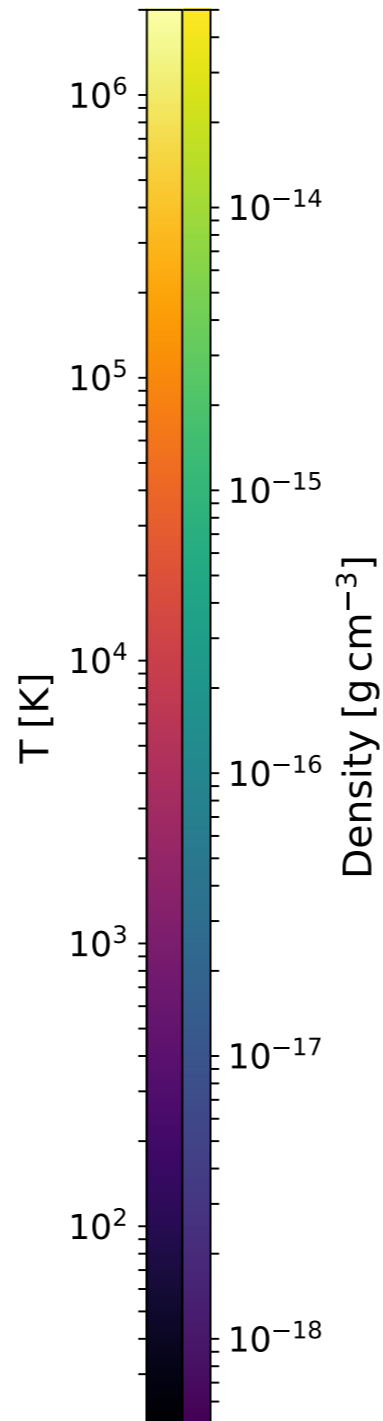
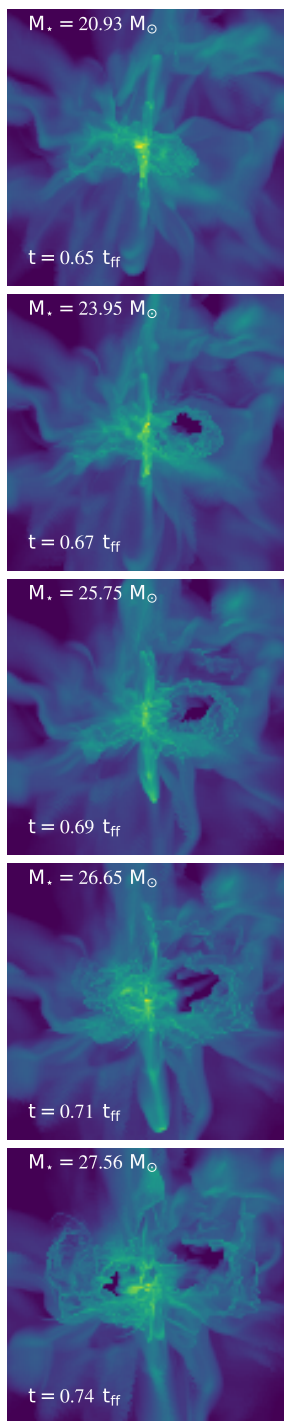
Radiation+Outflows+B

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Courtesy of Crystal Brogan

...**BUT WAIT!** What about fast, isotropic radiatively driven winds? **Caution: Preliminary results Stay tuned!**



$$\dot{M}_w \sim 10^{-7} - 10^{-6} M_\odot \text{ yr}^{-1}$$

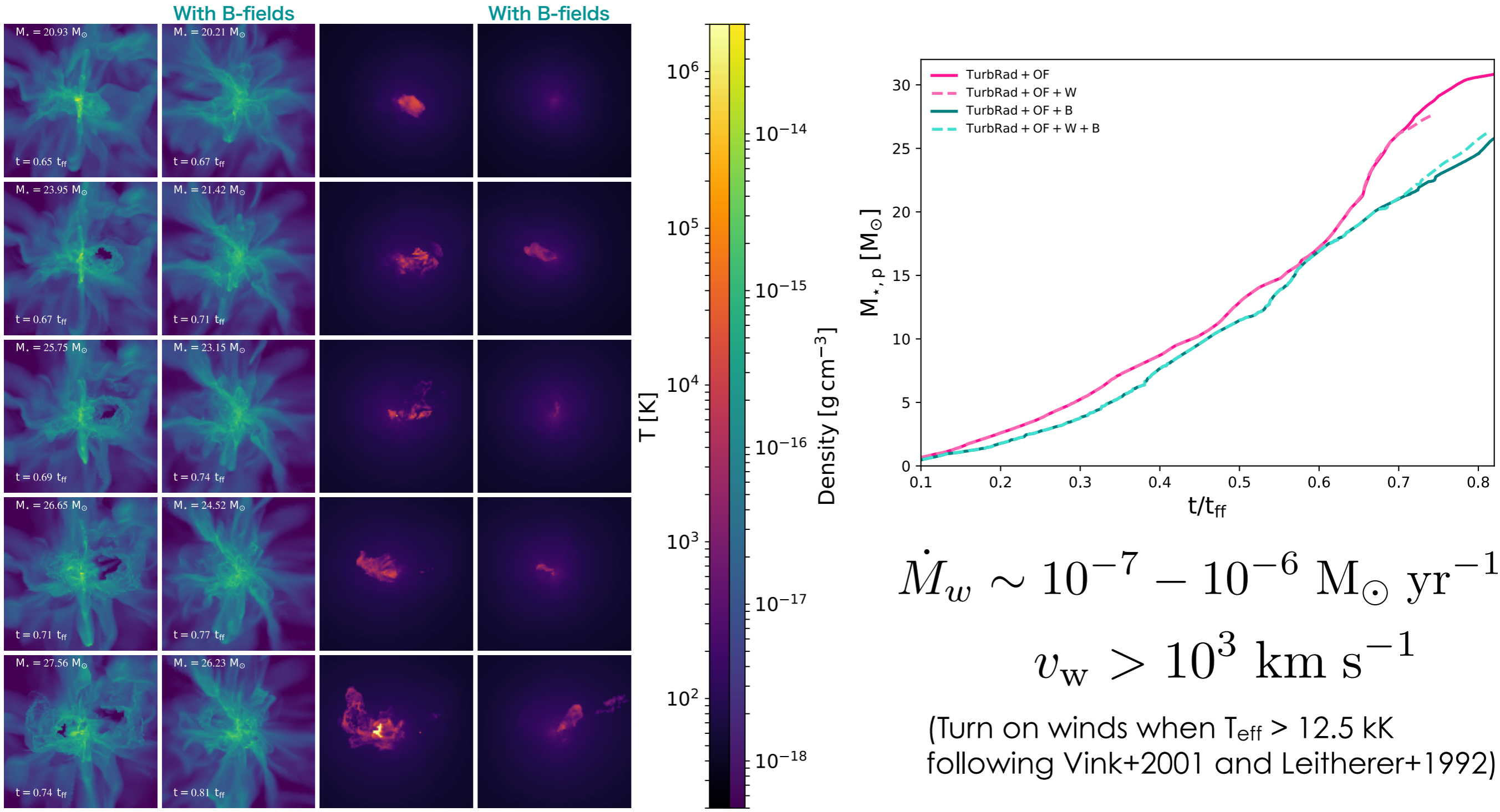
$$v_w > 10^3 \text{ km s}^{-1}$$

(Turn on winds when $T_{\text{eff}} > 12.5 \text{ kK}$
following Vink+2001 and Leitherer+1992)

(20,000 AU)²

Fast, Isotropic winds should **shock heat** gas yielding $T_{\text{gas}} \approx 10^6 \text{ K}$,
gas **adiabatically expands** reducing dM/dt

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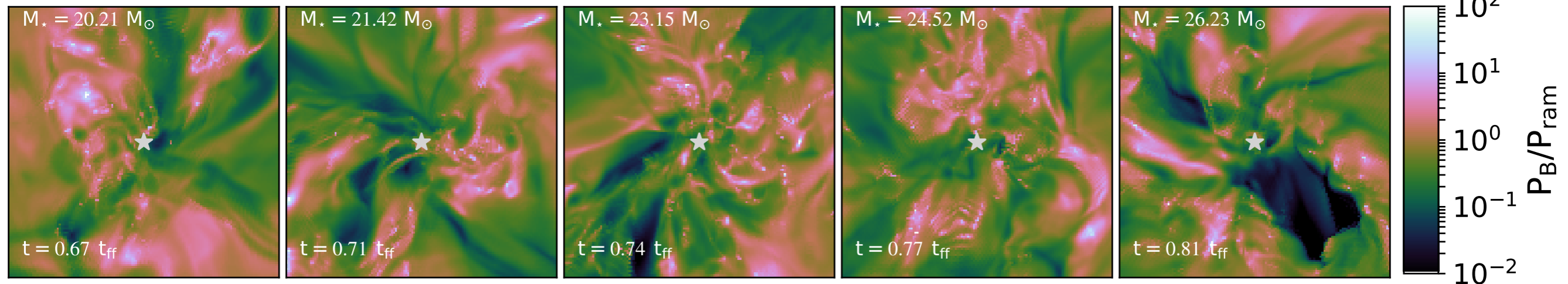
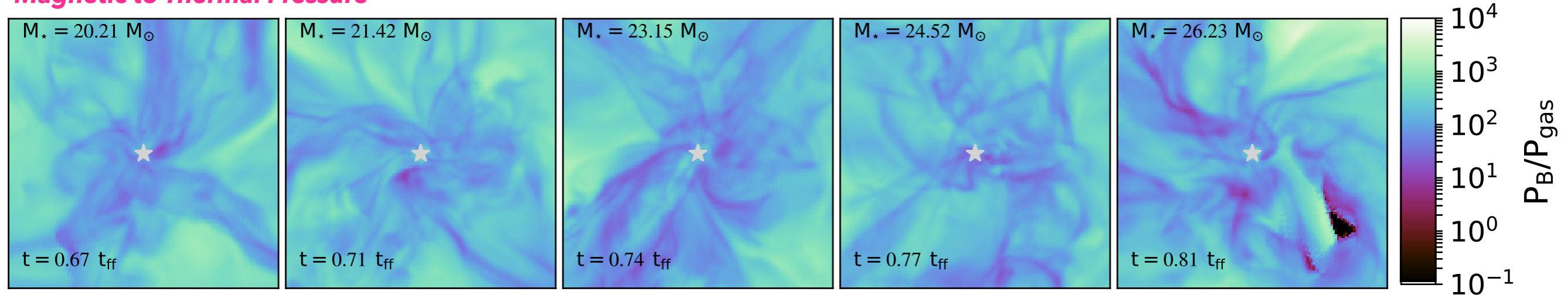
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Magnetic pressure **confines winds**, reducing shock heating and adiabatic expansion \rightarrow larger ρ and c_s such that dM/dt increases.

Magnetic to Thermal Pressure



Magnetic to Ram Pressure

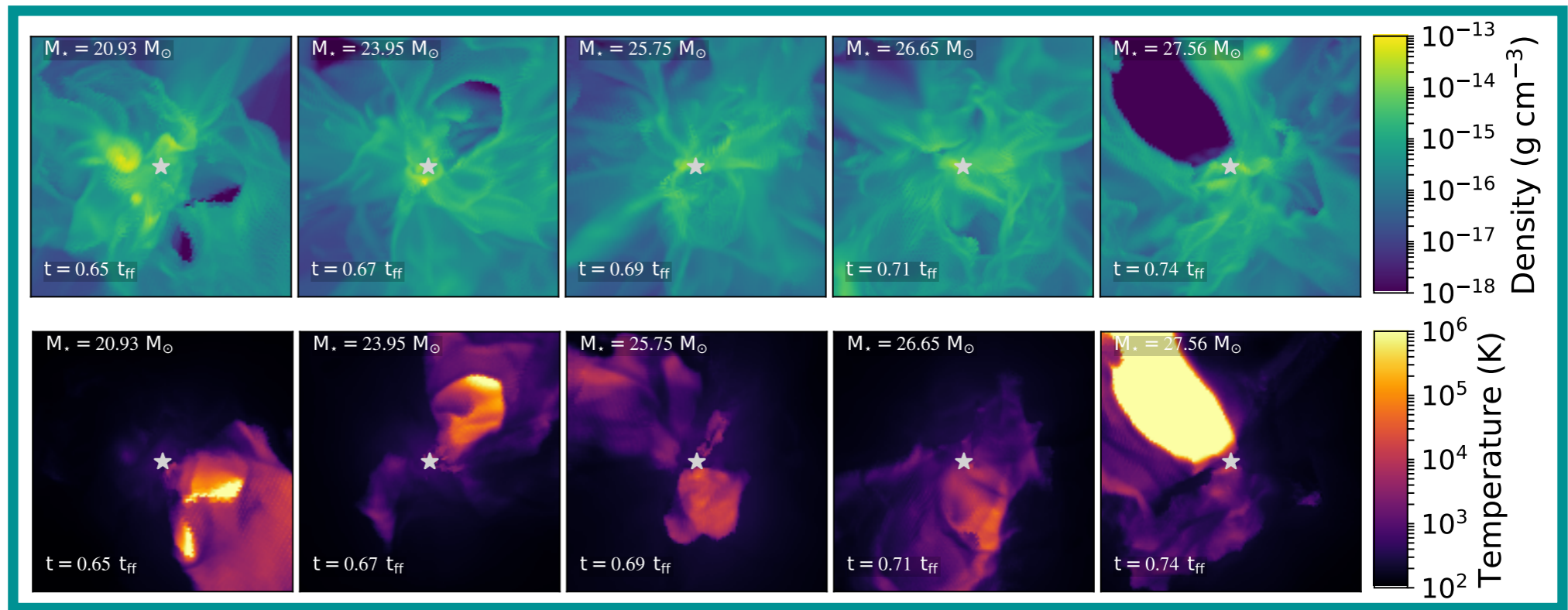
$(P_{\text{ram}} = \rho v^2)$

(4,000 AU)²

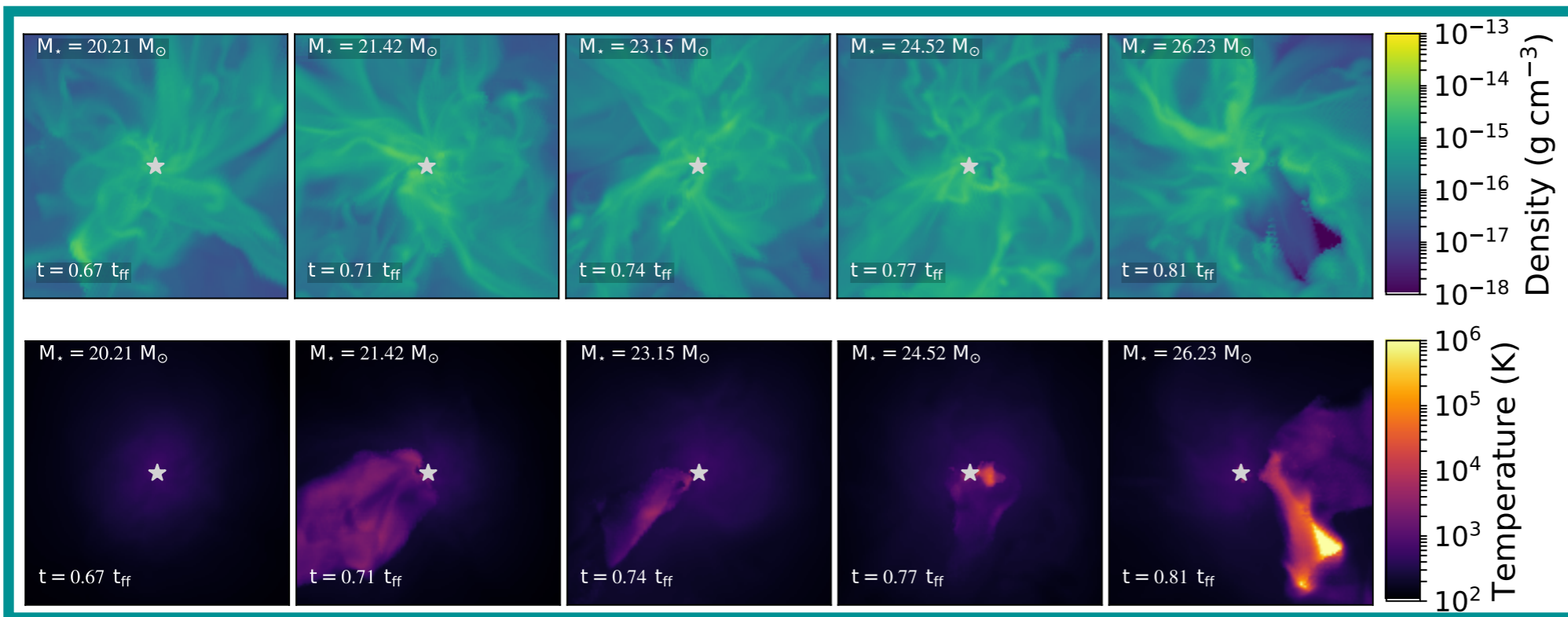
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Winds+ No B



Winds+B

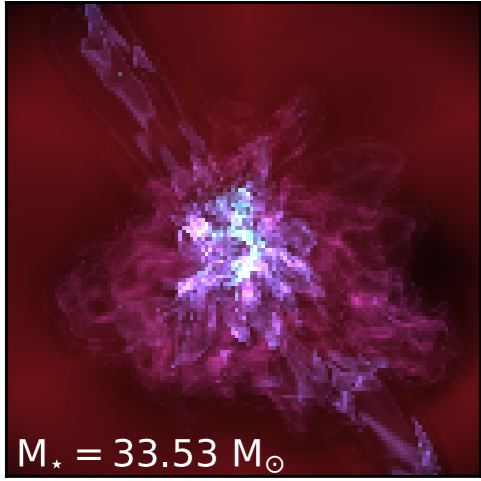


(4,000 AU)²

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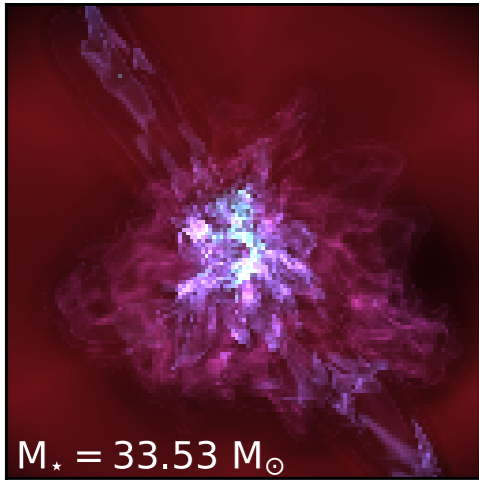
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Summary



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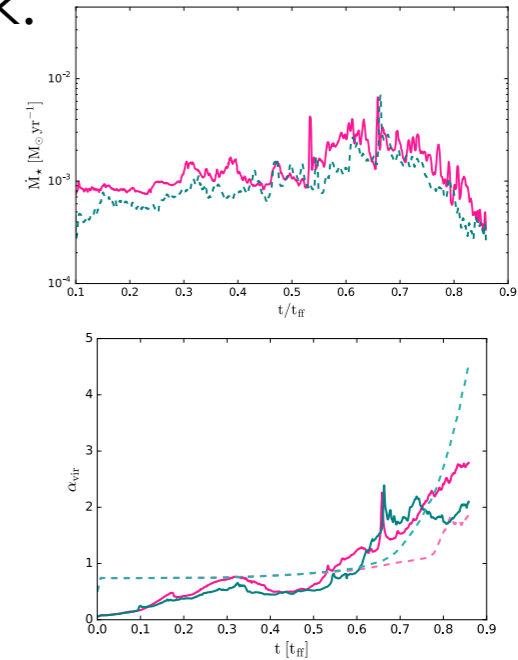
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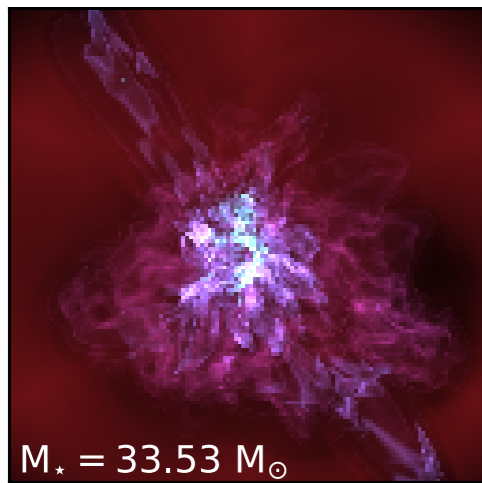
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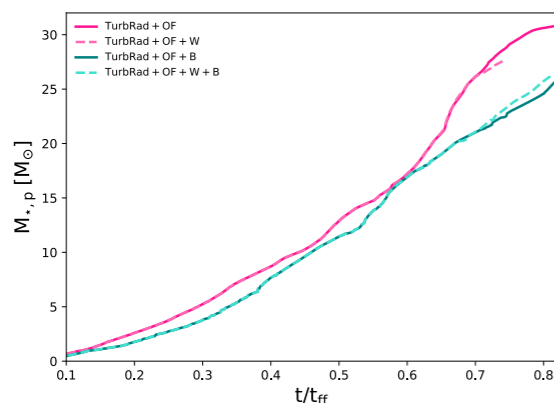
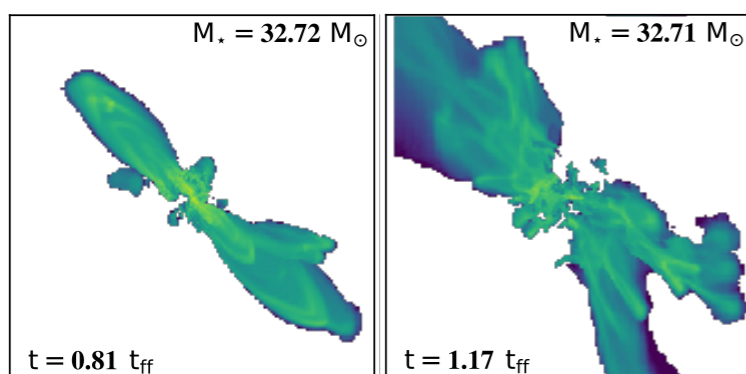
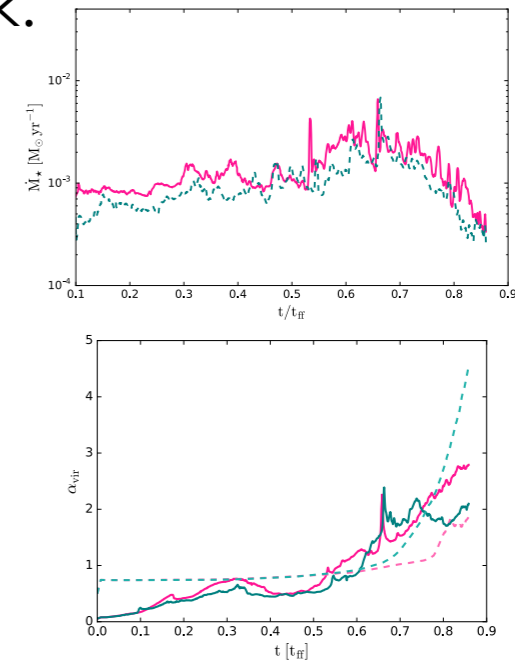
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Inclusion of magnetic fields in MSF:

- * Slows down the growth of massive stars
- * Significantly reduces formation of companions via turbulent fragmentation.
- * Leads to wider collimated molecular outflows.
- * When winds are included, leads to a positive (?) feedback effect (at least at early times)