## College of Science and Engineering School of Physics and Astronomy



## Radiation and Matter SCQF Level 11, U01439, PHY-4-RadMatt Thursday, 30th April 2009 9.30 - 11.30 a.m.

Chairman of Examiners

Professor J A Peacock

**External Examiner** 

Professor S Rawlings

## Answer **TWO** questions

The bracketed numbers give an indication of the value assigned to each portion of a question.

Only the supplied Electronic Calculators may be used during this examination.

Anonymity of the candidate will be maintained during the marking of this examination.

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1. State the Born-Oppenheimer approximation as it applies to small molecules.

State the rigid-rotor approximation for the rotational behaviour of the CO molecule and describe how it follows from the Born-Oppenheimer approximation.

[3]

[2]

The rigid-rotor Hamiltonian is

$$H = \frac{\mathbf{J}^2}{2I}$$

where **J** is the angular momentum operator and I the moment of inertia. Derive an expression for the reduced mass of the CO molecule. Taking the atomic nuclei of CO to be separated by R=1.1 Å on average, and the atomic masses to be integer multiples of the mass of the hydrogen atom, calculate the reduced mass of  $^{12}\mathrm{C}^{16}\mathrm{O}$  and the moment of inertia.

[4]

Given that the eigenvalues of  $J^2$  are  $J(J+1)\hbar^2$ , calculate the frequency of the  $J=1\to 0$  transition.

[5]

What are the wavefunctions appropriate to the rotational states?

[2]

Given that the permanent electric dipole moment is  $|e|\mathbf{d}$ , where e is the electronic charge and  $\mathbf{d}$  is a vector of length 0.023 Å, calculate the spontaneous emission rate of the  $J=1\to 0$  transition from the formula

$$w_{\text{spon}} = \frac{4\omega^3}{3\hbar c^3} \frac{e^2}{4\pi\epsilon_0} | < \text{upper } |\mathbf{r}| \text{ lower } > |^2$$

given that the lowest spherical harmonics are

[5]

$$Y_{0,0} = \sqrt{\frac{1}{4\pi}}, \quad Y_{1,\pm 1} = \mp \sqrt{\frac{3}{8\pi}} \sin \theta \ e^{\pm i\phi}, \quad Y_{1,0} = \sqrt{\frac{3}{4\pi}} \cos \theta.$$

Calculate the transition frequency and spontaneous emission rate of the  $J=1\to 0$  transition in the next most abundant isotopomer,  $^{13}\mathrm{C}^{16}\mathrm{O}$ .

[4]

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2. Given that in the wave zone the vector potential may be written in the general form

$$\mathbf{A}(\mathbf{r},t) = \mathbf{F}(t - \frac{\mathbf{r}.\mathbf{n}}{c})/R$$

where **F** is an arbitrary vector-valued function of a single argument (the retarded time), R (assumed constant) is the distance from the centre of the emitting region to the field point at **r**, and **n** is a unit vector in that same direction, derive expressions for  $\nabla \wedge \mathbf{F}$  and  $\nabla \cdot \mathbf{F}$  in terms of  $\frac{\partial \mathbf{F}}{\partial t}$ .

[2]

From Maxwell's equations in empty space

$$\nabla \wedge \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}, \quad \nabla \wedge \mathbf{B} = \epsilon_0 \mu_0 \frac{\partial \mathbf{E}}{\partial t}$$

and the fact that the scalar potential  $\phi = 0$  in the wave zone, show that  $\mathbf{n} \wedge \mathbf{E} = c\mathbf{B}$  and  $c\mathbf{n} \wedge \mathbf{B} = -\mathbf{E}$ . Infer that  $\mathbf{E}, \mathbf{B}$  and  $\mathbf{n}$  form a right-handed triad like x, y and z.

[2]

Define the Poynting vector S.

[2] [2]

Show that

$$\mathbf{S} = \epsilon_0 c \mathbf{\dot{A}}_{\perp}^2 \mathbf{n}$$

[4]

Given the Liénard-Wiechert potential in the wave zone for radiation from a single particle of charge q and velocity  $\mathbf{v}$  in the non-relativistic limit

$$\mathbf{A}(\mathbf{r},t) = \frac{\mu_0 q \mathbf{v}}{4\pi R}$$

show that

$$\mathbf{S} = \frac{1}{4\pi R^2} \frac{q^2}{4\pi c^3 \epsilon_0} \dot{v}^2 \sin^2 \alpha \, \mathbf{n}$$

[4]

where  $\alpha$  is the polar angle. Finally, integrate this over the sphere to obtain Larmor's formula for the power radiated by an accelerated charge

$$P = \frac{q^2 \dot{v}^2}{6\pi c^3 \epsilon_0}$$

[3]

A particle of mass m and charge Q initially moves with a velocity  $\mathbf{v}$  perpendicular to a uniform magnetic field of strength  $\mathbf{B}_0$ . Show that the trajectory is a circle, and derive the radius of the circle r and the angular frequency  $\omega$ .

[2]

Show that the particle radiates electromagnetic waves at a rate

$$P = \frac{q^4}{6\pi c^3 \epsilon_0} \left(\frac{vB_0}{m}\right)^2$$

[2]

and that the resulting loss of kinetic energy causes the velocity to decrease as

$$v = v_0 \exp(-t/T)$$

giving an expression for T.

[2]

## Radiation and Matter (U01439)

Define the specific intensity  $I_{\nu}$  and give its (MKS) units. [3]Which two aspects of the radiation field does  $I_{\nu}$  not encapsulate? [2]The equation of radiative transfer in standard notation is  $\frac{dI_{\nu}}{dl} = -\kappa_{\nu}I_{\nu} + \epsilon_{\nu}.$ What are  $\kappa_{\nu}$  and  $\epsilon_{\nu}$  called, and what are their units? [4]Show that  $I_{\nu}$  is constant in empty space. [1]It follows that  $I_{\nu}$  measured by an observer looking at the centre of the solar disc is independent of his distance from the Sun; how is this consistent with the inverse-square law? [2] $\kappa_{\nu}$  and  $\epsilon_{\nu}$  for any spectral line in any astrophysical setting both depend on the frequency as some function  $\phi(\nu)$ . How is this usually normalized? [1]Describe all the processes you can think of that can contribute to the form of  $\phi(\nu)$ in various astronomical environments, indicating where appropriate the standard form  $\phi(\nu)$  might take if that process were acting alone. [5][1]Show that these processes fall into two main classes. A uniform interstellar cloud of thickness L along the line of sight absorbs and emits radiation with constant  $\kappa_{\nu}$  and  $\epsilon_{\nu}$  at a certain frequency  $\nu$ . Solve for the observed  $I_{\nu}$  assuming that no light is incident on the far side of the cloud. [3]How does the solution behave in the limits: a)  $L\kappa_{\nu} \ll 1$  and b)  $L\kappa_{\nu} \gg 1$ ? [2]

How would b) be affected if there were a uniform source on the far side of the

[1]

cloud?